

Lecture 10 - Topics

- Relativistic strings: Nambu-Gotto action, equations of motion and boundary conditions

Reading: Section 6.1 - 6.5

$S \propto$ world-sheet area (A)

Units of const. of proportionality

$$[S] = [\int dtL] = \text{Energy} \times \text{time} = \frac{ML^2}{T}$$

$$[\text{World Sheet Area}] = L^2$$

$$[\text{Const. of Proportionality}] = \frac{M}{T}$$

Recall for tension T_0 : $[T_0] = [\text{Force}] = \frac{ML}{T^2} \cdot \left[\frac{T_0}{c} \right] = \frac{M}{T}$. So let T_0/c be constant of proportionality for $S \propto A$.

Nambu-Gotto Action

$$S = -\frac{T_0}{c} \int_{\tau_i}^{\tau_f} d\tau \int_0^{\sigma_1} d\sigma \sqrt{(\dot{X} \cdot X')^2 - (\dot{X})^2 (X')^2}$$

where $\dot{X}^\mu = \frac{\partial X^\mu}{\partial \tau}$. $X' = \frac{\partial X^\mu}{\partial \sigma}$. The two tangents to the surface.

This defines the dynamics of the string.

World sheet metric: Let $\zeta^1 = \tau$, $\zeta^2 = \sigma$.

Distance between 2 points on the worldsheet:

$$\begin{aligned} -dx^2 &= dX^\mu dX_\mu = \eta_{\mu\nu} dX^\mu dX^\nu \\ &= \eta_{\mu\nu} \frac{\partial X^\mu}{\partial \zeta^\alpha} \frac{\partial X^\nu}{\partial \zeta^\beta} d\zeta^\alpha d\zeta^\beta \\ &= \gamma_{\alpha\beta}(\zeta) d\zeta^\alpha d\zeta^\beta \end{aligned}$$

where $\gamma_{\alpha\beta} = \eta_{\mu\nu} \frac{\partial X^\mu}{\partial \zeta^\alpha} \frac{\partial X^\nu}{\partial \zeta^\beta}$

$$\gamma_{\alpha\beta} = \begin{pmatrix} \dot{X}^2 & \dot{X}X' \\ \dot{X}X' & (X')^2 \end{pmatrix}$$

String world-sheet metric. All signs from the η are in the X 's. So:

$$S = -\frac{T_0}{c} \iint d\tau d\sigma \sqrt{-\gamma}$$

where $\gamma = \det(\gamma_{\alpha\beta})$. Lorentz invar., reparam invar.

$$X = \int_{\tau_i}^{\tau_f} d\tau \int_0^{\sigma_1} d\sigma \mathcal{L}(\dot{X}^\mu, X'^\mu)$$

Lagrangian Density:

$$\mathcal{L}(\dot{X}^\mu, X'^\mu) = -\frac{T_0}{c} \sqrt{(\dot{X} \cdot X')^2 - (\dot{X})^2 (X')^2}$$

$$\delta S = 0$$

$$\delta S = \int_{\tau_i}^{\tau_f} d\tau \int_0^{\sigma_1} d\sigma \left[\left(\frac{\partial \mathcal{L}}{\partial \dot{X}^\mu} \right) \left(\frac{\partial(\delta X^\mu)}{\partial \tau} \right) + \left(\frac{\partial \mathcal{L}}{\partial X'^\mu} \right) \left(\frac{\partial(\delta X^\mu)}{\partial \sigma} \right) \right]$$

$$\delta \dot{X}^\mu = \delta \frac{\partial X^\mu}{\partial \tau}$$

Define:

$$\frac{\partial \mathcal{L}}{\partial \dot{X}^\mu} = \mathcal{P}_\mu^\tau = -\frac{T_0}{c} \frac{(\dot{X} \cdot X') X'_\mu - (X')^2 \dot{X}_\mu}{\sqrt{(\dot{X} \cdot X')^2 - (\dot{X})^2 (X')^2}}$$

$$\frac{\partial \mathcal{L}}{\partial X'^\mu} = \mathcal{P}_\mu^\sigma = -\frac{T_0}{c} \frac{(\dot{X} X') \dot{X}_\mu - (\dot{X})^2 X'_\mu}{\sqrt{\dots}}$$

Things will simplify soon.

So:

$$\delta S = \iint d\tau d\sigma \left[\frac{\partial}{\partial \tau} (\delta X^\mu \mathcal{P}_\mu^\tau) + \frac{\partial}{\partial \sigma} (\delta X^\mu \mathcal{P}_\mu^\sigma) - \delta X^\mu \left(\frac{\partial \mathcal{P}_\mu^\tau}{\partial \tau} + \frac{\partial \mathcal{P}_\mu^\sigma}{\partial \sigma} \right) \right]$$

$\delta S = 0$ so ...

$$\frac{\partial \mathcal{P}_\mu^\tau}{\partial \tau} + \frac{\partial \mathcal{P}_\mu^\sigma}{\partial \sigma} = 0$$

For open strings:

$$\int_{\tau_i}^{\tau_f} d\tau [\delta X^0(\tau, \sigma_1) \mathcal{P}_0^\sigma(\tau, \sigma_1) - \delta X^0(\tau, 0) \mathcal{P}_0^\sigma(\tau, \sigma_1) + \delta X^1(\tau, \sigma_1) \mathcal{P}_1^\sigma(\tau, \sigma_1) - \delta X^1(\tau, 0) \mathcal{P}_1^\sigma(\tau, \sigma_1) + \dots + \delta X^d \dots]$$

2D Constraints.

For most, get choice as to how the term vanishes since product of 2 terms (so either can be 0).

For $\mu \neq 0$:

1. Dirichlet BCs: $X^\mu(\tau, \sigma_*) = \text{constant}$, $\delta X^\mu(\tau, \sigma_*) = 0$ (for $\sigma_* = 0$ or σ_1)
2. Free BCs: $P_\mu^\sigma(\tau, \sigma_*) = 0$ for $\sigma_* = 0$ or σ_1 .

For $\mu = 0$:

1. $\partial X^0 / \partial \tau < 0$ can't have Dirichlet.
2. Free BCs. $\mathcal{P}_0^\sigma(\tau, 0) = \mathcal{P}_0^\tau(\tau, \sigma_1) = 0$.