

# \* Radiative Properties of Gases

o-n-o Molecules

vibrati<sup>n</sup>

$$E_n = \hbar \omega_0 (n + \frac{1}{2})$$

rotation

$$E_l = \frac{\hbar^2}{2I} l(l+1) \quad \left\{ \begin{array}{l} l=0, 1, \dots \\ m \leq |l| \end{array} \right.$$

$$\hbar \omega_p = E_{n+1} - E_n \quad (\text{select rule})$$

$$\omega_p = \omega_0$$

$$\eta = \frac{2\pi l}{\lambda} = \frac{\omega}{c}$$

$$\hbar \omega_p = E_{l+1} - E_l$$

$$\omega_p = \frac{\hbar}{2I} [(l+1)(l+2) - l(l+1)] = \frac{\hbar}{I} (l+1)$$

~~l depends on~~

vibrational - rotational

$$E_{nl} = n\hbar(n + \frac{1}{2}) + \frac{\hbar^2}{2I_n} l(l+1)$$

$$\omega_p \text{ or } \eta_p = \frac{\omega_p}{c}$$

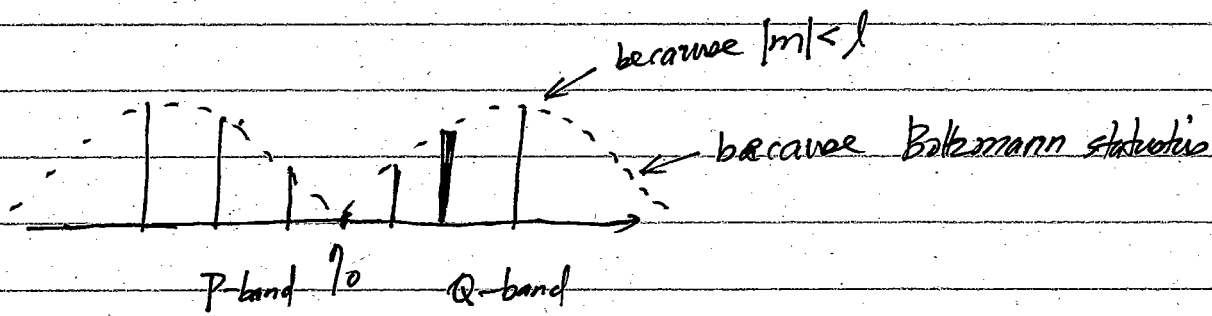
$$\eta_p = \eta_0 - (B_{n+1} + B_n)l + (B_{n+1} - B_n)l^2 \quad \Delta l = -1$$

j=1, 2

$$\eta_Q = \eta_0 + (B_{n+1} - B_n)l + (B_{n+1} - B_n)l^2 \quad \Delta l = 0$$

$$\eta_R = \eta_0 + 2B_{n+1} + (3B_{n+1} - B_n)l + (B_{n+1} - B_n)l^2 \quad \Delta l = 1$$

↑  
R-band



Complications: ① for a single molecules  $\text{CO}_2$   
 there can be different  $\eta_0$   
 ② Each line can be broadened.

a) natural energy uncertainty

Heisenberg uncertainty  $\Delta E \cdot \Delta t \geq \frac{\hbar}{2}$

b) molecules moving: Doppler effect

$$\eta_{\text{obs}} = \eta_{\text{em}} \left( 1 + \frac{\vec{v} \cdot \vec{s}}{c} \right)$$

$$p(v) = \left( \frac{m}{2\pi kT} \right)^{\frac{3}{2}} 4\pi v^2 \exp\left(-\frac{mv^2}{2kT}\right)$$

$$\Rightarrow K_{\eta} = \sqrt{\frac{\ln 2}{\pi}} \left( \frac{s}{b_0} \right) \exp\left[-(\ln 2) \left( \frac{\eta - \eta_0}{b_0} \right)^2\right]$$

$$b_0 = \frac{\eta_0}{c_0} \sqrt{\frac{2kT}{m \ln 2}}$$

c) collision

$$K_{\eta} = \frac{s}{\pi} \frac{b_c}{(\eta - \eta_0)^2 + b_c^2}$$

$$s = \int s_{\eta} K_{\eta} d\eta$$

$$b_c = b_{c0} \left( \frac{p}{p_0} \right) \sqrt{\frac{T_0}{T}}$$

Recall Harmonic Oscillator under an electric field

$$\Delta x = x_0 e^{-i\omega t}$$

$$x_0 = \frac{E_0/m}{-\omega^2 + \omega_0^2 + i\gamma\omega}, \quad \Sigma = 1 + \frac{Ne^2/m}{-\omega^2 + \omega_0^2 + i\gamma\omega}$$

For solid, we sum  $N$  oscillators to talk about  $\vec{P} \propto \Sigma \mathcal{E}$

For gases,  $N \approx 1$   $N$  small compared to solid.  
imaginary part of interest

$$\Sigma'' = \frac{Ne^2/m\epsilon_0 (\gamma\omega)}{(\omega_0^2 - \omega^2)^2 + (\gamma\omega)^2}$$

$$2nk = \Sigma''$$

$$k = \frac{1}{2} \frac{Ne^2}{m\epsilon_0} \frac{\gamma\omega}{(\omega^2 - \omega_0^2)^2 + (\gamma\omega)^2}$$

$$\approx \frac{1}{2} \frac{Ne^2}{m\epsilon_0} \frac{\gamma\omega}{(\omega - \omega_0)^2 + 4\omega^2 + (\gamma\omega)^2}$$

absorption coefficient

$$\alpha_j = \frac{4\pi k_j}{\lambda} = \frac{2\pi Ne^2}{m\epsilon_0 c_0} \frac{\gamma_j}{4(\omega - \omega_0)^2 + \gamma_j^2}$$

back use  $k_j \propto \omega$   $\hookrightarrow$  Lorentzian profile

EM theory.  $\frac{dI_\lambda}{d\lambda} = -\alpha_\lambda I_\lambda$

$$K_\eta = \frac{S}{\pi} \frac{b_c}{(\eta - \eta_0)^2 + b_c^2}$$

$\hookrightarrow b_c$  - line half width

$$S \equiv \int_{\lambda_1}^{\lambda_2} K_\eta d\eta$$

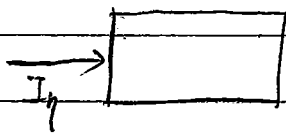
$\uparrow$  line strength (line integrated absorption)

We can also define mass absorpt. coefficient

$$K_{a\eta} = \frac{k_{\eta a}}{\rho_a}$$

$\rho_a$  — partial density of species a.

$S$  — optical path length  
 $\rho_a S$  — mass based



$$\frac{dI_\eta}{ds} = -k_\eta I_\eta$$

$$I_\eta = I_\eta(0) e^{-k_\eta s}$$

absorbance over  $S$

$$A_\eta = \frac{I_\eta(0) - I_\eta(s)}{I_\eta(0)} = 1 - e^{-k_\eta X}$$

$X = s$   
 $\rho_a s$  — mass path length

$$\text{Kirchhoff law} = \epsilon_\eta$$

Integrated line Emissivity

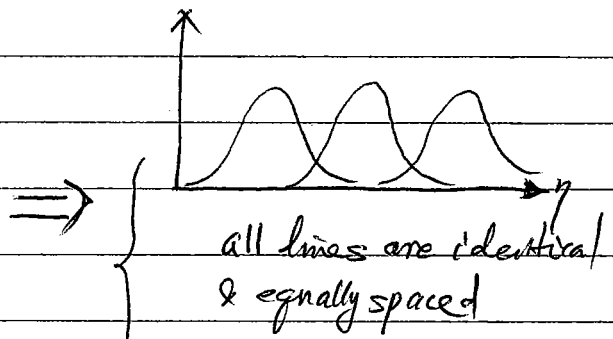
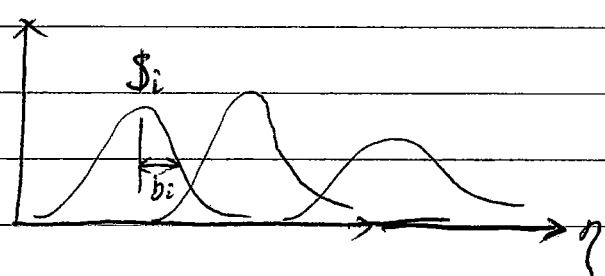
$$W = \frac{\int_0^\infty \epsilon_\eta I_{b\eta} d\eta}{\int_0^\infty I_{b\eta} d\eta}$$

$\hookrightarrow$  blackbody varies little over one line width

$$= \int_0^\infty (1 - e^{-k_\eta X}) d\eta$$

$$= \begin{cases} SX & (kX \ll 1) \text{ weak line} \\ 2\sqrt{SXb} & (kX \gg 1) \text{ strong line} \end{cases}$$

How to deal with overlapping lines



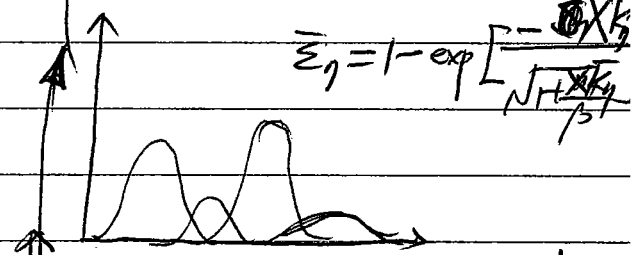
all lines are identical & equally spaced

Elsasser Model

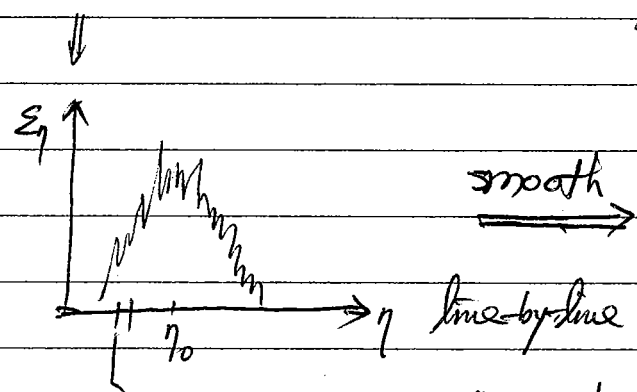
$$K_\eta = \sum_j K_{\eta_j}$$

$$\epsilon_\eta = 1 - \exp \left[ -X \sum_j K_{\eta_j} \right]$$

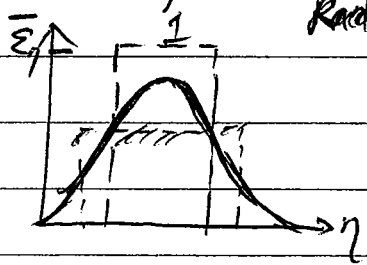
$$\bar{\epsilon}_\eta = 1 - \exp \left[ -\frac{X \sum_j K_{\eta_j}}{\sqrt{\frac{X \sum_j K_{\eta_j}}{N}}} \right]$$



Goody model - statistical Randomly spaced



smooth



$$\bar{\epsilon} = \frac{1}{2\eta} \int_{-\infty}^{\infty} \epsilon_\eta d\eta$$

The above are narrow band models.

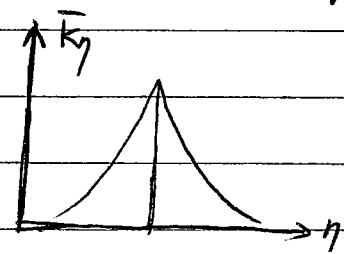
Wide Band Model.

$$A \equiv \int_{\text{band}} \bar{\epsilon}_\eta d\eta$$

effective band width if \$\epsilon=1\$

$$\bar{\epsilon}_\eta = 1 - e^{-\bar{K}_\eta X}$$

$$\bar{K}_\eta = \frac{\sum_i K_i}{d}$$



Construct models for \$\bar{K}\_\eta\$ rather than \$\bar{\epsilon}\_\eta\$

So that no X-dependence

Edwards exponential.