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2.161 Signal Processing: Continuous and Discrete  
Fall 2008

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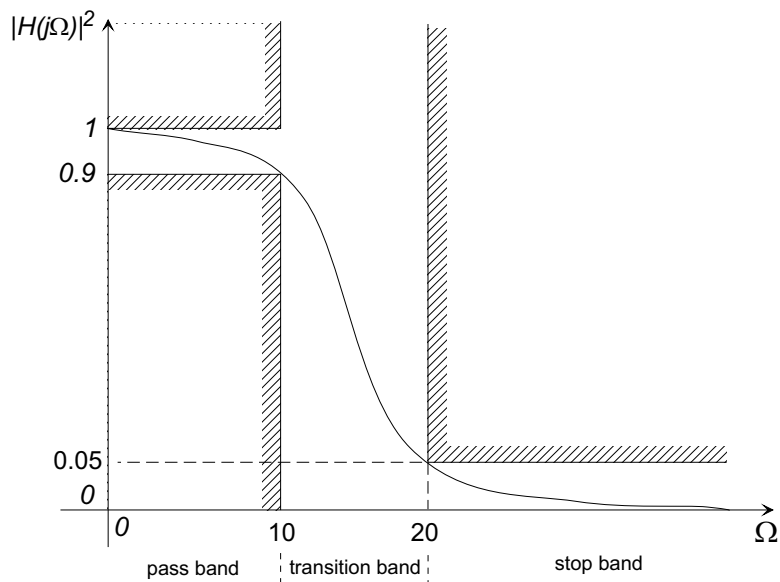
**Lecture 7**<sup>1</sup>

**Reading:**

- Class handout: *Introduction to Continuous Time Filter Design.*

**1 Butterworth Filter Design Example**

(Same problem as in the Class Handout). Design a Butterworth low-pass filter to meet the power gain specifications shown below:



At the two critical frequencies

$$\frac{1}{1 + \epsilon^2} = 0.9 \quad \longrightarrow \quad \epsilon = 0.3333$$
$$\frac{1}{1 + \lambda^2} = 0.05 \quad \longrightarrow \quad \lambda = 4.358$$

Then

$$N \geq \frac{\log(\lambda/\epsilon)}{\log(\Omega_r/\Omega_c)} = 3.70$$

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we therefore select  $N=4$ . The 4 poles ( $p_1 \dots p_4$ ) lie on a circle of radius  $r = \Omega_c \epsilon^{-1/N} = 13.16$  and are given by

$$\begin{aligned} |p_n| &= 13.16 \\ \angle p_n &= \pi(2n + 3)/8 \end{aligned}$$

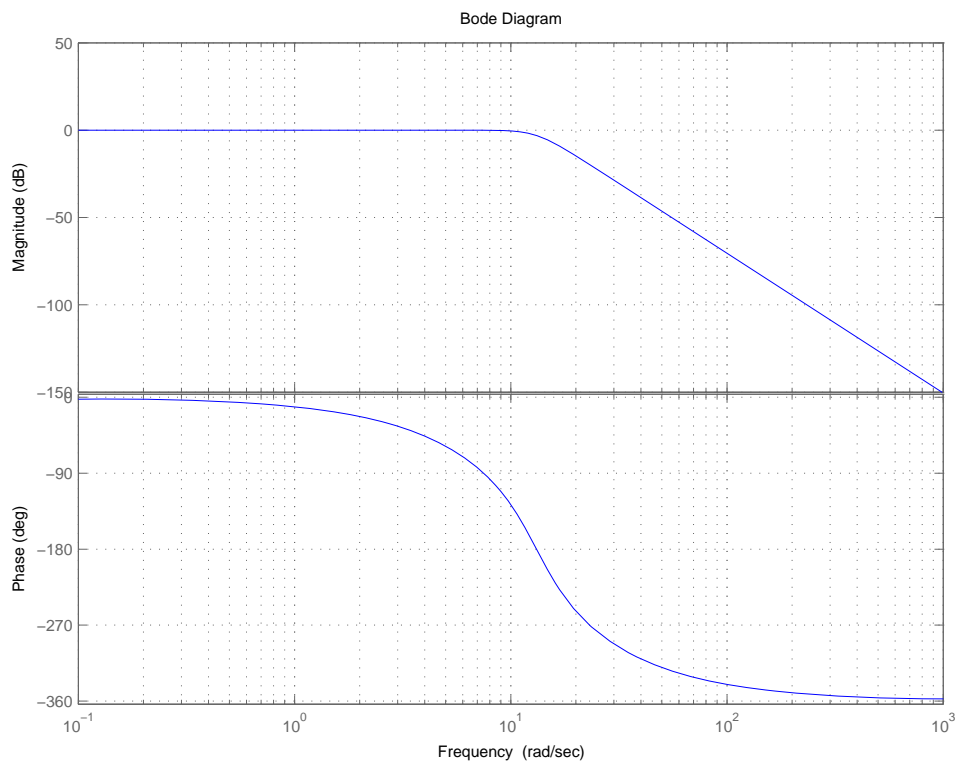
for  $n = 1 \dots 4$ , giving a pair of complex conjugate pole pairs

$$\begin{aligned} p_{1,4} &= -5.04 \pm j12.16 \\ p_{2,3} &= -12.16 \pm j5.04 \end{aligned}$$

The transfer function, normalized to unity gain, is

$$H(s) = \frac{29993}{(s^2 + 10.07s + 173.2)(s^2 + 24.32s + 173.2)}$$

and the filter Bode plots are shown below.



## 2 Chebyshev Filters

The order of a filter required to meet a low-pass specification may often be reduced by relaxing the requirement of a monotonically decreasing power gain with frequency, and allowing

“ripple” to occur in either the pass-band or the stop-band. The Chebyshev filters allow these conditions:

$$\text{Type 1} \quad |H(j\Omega)|^2 = \frac{1}{1 + \epsilon^2 T_N^2(\Omega/\Omega_c)} \quad (1)$$

$$\text{Type 2} \quad |H(j\Omega)|^2 = \frac{1}{1 + \epsilon^2 (T_N^2(\Omega_r/\Omega_c)/T_N^2(\Omega_r/\Omega))} \quad (2)$$

Where  $T_N(x)$  is the Chebyshev polynomial of degree  $N$ . Note the similarity of the form of the Type 1 power gain (Eq. (1)) to that of the Butterworth filter, where the function  $T_N(\Omega/\Omega_c)$  has replaced  $(\Omega/\Omega_c)^N$ . The Type 1 filter produces an all-pole design with slightly different pole placement from the Butterworth filters, allowing resonant peaks in the pass-band to introduce ripple, while the Type 2 filter introduces a set of zeros on the imaginary axis above  $\Omega_r$ , causing a ripple in the stop-band.

The Chebyshev polynomials are defined recursively as follows

$$\begin{aligned} T_0(x) &= 1 \\ T_1(x) &= x \\ T_2(x) &= 2x^2 - 1 \\ T_3(x) &= 4x^3 - 3x \\ &\vdots \\ T_N(x) &= 2xT_{N-1}(x) - T_{N-2}(x), \quad N > 1 \end{aligned} \quad (3)$$

with alternate definitions

$$T_N(x) = \cos(N \cos^{-1}(x)) \quad (4)$$

$$= \cosh(N \cosh^{-1}(x)) \quad (5)$$

The Chebyshev polynomials have the *min-max* property:

*Of all polynomials of degree  $N$  with leading coefficient equal to one, the polynomial*

$$T_N(x)/2^{N-1}$$

*has the smallest magnitude in the interval  $|x| \leq 1$ . This “minimum maximum” amplitude is  $2^{1-N}$ .*

In low-pass filters given by Eqs. (13) and (14), this property translates to the following characteristics:

<b>Filter</b>	<b>Pass-Band Characteristic</b>	<b>Stop-Band Characteristic</b>
Butterworth	Maximally flat	Maximally flat
Chebyshev Type 1	Ripple between 1 and $1/(1 + \epsilon^2)$	Maximally flat
Chebyshev Type 2	Maximally flat	Ripple between 1 and $1/(1 + \lambda^2)$

## 2.1 The Chebyshev Type 1 Filter

With the power response from Eq. (13)

$$|H(j\Omega)|^2 = \frac{1}{1 + \epsilon^2 T_N^2(\Omega/\Omega_c)}$$

and the filter specification from Fig. 1, the required filter order may be found as follows. At the edge of the stop-band  $\Omega = \Omega_r$

$$|H(j\Omega_r)|^2 = \frac{1}{1 + \epsilon^2 T_N^2(\Omega_r/\Omega_c)} \leq \frac{1}{1 + \lambda^2}$$

so that

$$\lambda \leq \epsilon T_N(\Omega_r/\Omega_c) = \epsilon \cosh(N \cosh^{-1}(\Omega_r/\Omega_c))$$

and solving for  $N$

$$N \geq \frac{\cosh^{-1}(\lambda/\epsilon)}{\cosh^{-1}(\Omega_r/\Omega_c)} \quad (6)$$

The characteristic equation of the power transfer function is

$$1 + \epsilon^2 T_N^2\left(\frac{s}{j\Omega_c}\right) = 0 \quad \text{or} \quad T_N\left(\frac{s}{j\Omega_c}\right) = \pm \frac{j}{\epsilon}$$

Now  $T_N(x) = \cos(N \cos^{-1}(x))$ , so that

$$\cos\left(N \cos^{-1}\left(\frac{s}{j\Omega_c}\right)\right) = \pm \frac{j}{\epsilon} \quad (7)$$

If we write  $\cos^{-1}(s/j\Omega_c) = \gamma + j\alpha$ , then

$$\begin{aligned} s &= \Omega_c (j \cos(\gamma + j\alpha)) \\ &= \Omega_c (\sinh \alpha \sin \gamma + j \cosh \alpha \cos \gamma) \end{aligned} \quad (8)$$

which defines an ellipse of width  $2\Omega_c \sinh(\alpha)$  and height  $2\Omega_c \cosh(\alpha)$  in the  $s$ -plane. The poles will lie on this ellipse. Substituting into Eq. (16)

$$\begin{aligned} T_N\left(\frac{s}{j\Omega_c}\right) &= \cos(N(\gamma + j\alpha)) \\ &= \cos N\gamma \cosh N\alpha - j \sin N\gamma \sinh N\alpha, \end{aligned}$$

the characteristic equation becomes

$$\cos N\gamma \cosh N\alpha - j \sin N\gamma \sinh N\alpha = \pm \frac{j}{\epsilon}. \quad (9)$$

Equating the real and imaginary parts in Eq. (21), (1) since  $\cosh x \neq 0$  for real  $x$  we require  $\cos N\gamma = 0$ , or

$$\gamma_n = \frac{(2n-1)\pi}{2N} \quad n = 1, \dots, 2N \quad (10)$$

and, (2) since at these values of  $\gamma$ ,  $\sin N\gamma = \pm 1$  we have

$$\alpha = \pm \frac{1}{N} \sinh^{-1} \frac{1}{\epsilon} \quad (11)$$

As in the Butterworth design procedure, we select the left half-plane poles as the poles of the filter frequency response.

**Design Procedure:**

1. Determine the filter order

$$N \geq \frac{\cosh^{-1}(\lambda/\epsilon)}{\cosh^{-1}(\Omega_r/\Omega_c)}$$

2. Determine  $\alpha$

$$\alpha = \pm \frac{1}{N} \sinh^{-1} \frac{1}{\epsilon}$$

3. Determine  $\gamma_n$ ,  $n = 1 \dots N$

$$\gamma_n = \frac{(2n-1)\pi}{2N} \quad n = 1, \dots, N$$

4. Determine the  $N$  left half-plane poles

$$p_n = \Omega_c (\sinh \alpha \sin \gamma_n + j \cosh \alpha \cos \gamma_n) \quad n = 1, \dots, N$$

5. Form the transfer function

- (a) If  $N$  is odd

$$H(s) = \frac{-p_1 p_2 \dots p_N}{(s - p_1)(s - p_2) \dots (s - p_N)}$$

- (b) If  $N$  is even

$$H(s) = \frac{1}{1 + \epsilon^2} \frac{p_1 p_2 \dots p_N}{(s - p_1)(s - p_2) \dots (s - p_N)}$$

The difference in the gain constants in the two cases arises because of the ripple in the pass-band. When  $N$  is odd, the response  $|H(j0)|^2 = 1$ , whereas if  $N$  is even the value of  $|H(j0)|^2 = 1/(1 + \epsilon^2)$ .

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### ■ Example 1

Repeat the previous Butterworth design example using a Chebyshev Type 1 design.

From the previous example we have  $\Omega_c = 10$  rad/s.,  $\Omega_r = 20$  rad/s.,  $\epsilon = 0.3333$ ,  $\lambda = 4.358$ . The required order is

$$N \geq \frac{\cosh^{-1}(\lambda/\epsilon)}{\cosh^{-1}(\Omega_r/\Omega_c)} = \frac{\cosh^{-1} 13.07}{\cosh^{-1} 2} = 2.47$$

Therefore take  $N = 3$ . Determine  $\alpha$ :

$$\alpha = \frac{1}{N} \sinh^{-1} \left( \frac{1}{\epsilon} \right) = \frac{1}{3} \sinh^{-1}(3) = 0.6061$$

and  $\sinh \alpha = 0.6438$ , and  $\cosh \alpha = 1.189$ . Also,  $\gamma_n = (2n - 1)\pi/6$  for  $n = 1 \dots 6$  as follows:

$n:$	1	2	3	4	5	6
$\gamma_n:$	$\pi/6$	$\pi/2$	$5\pi/6$	$7\pi/6$	$3\pi/2$	$11\pi/6$
$\sin \gamma_n:$	$1/2$	1	$1/2$	$-1/2$	-1	$-1/2$
$\cos \gamma_n:$	$\sqrt{3}/2$	0	$-\sqrt{3}/2$	$-\sqrt{3}/2$	0	$\sqrt{3}/2$

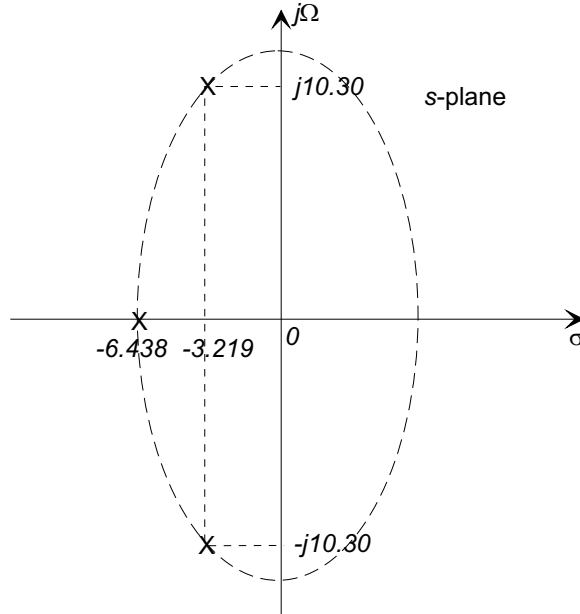
Then the poles are

$$\begin{aligned}
 p_n &= \Omega_c (\sinh \alpha \sin \gamma_n + j \cosh \alpha \cos \gamma_n) \\
 p_1 &= 10 \left( 0.6438 \times \frac{1}{2} + j1.189 \times \frac{\sqrt{3}}{2} \right) = 3.219 + j10.30 \\
 p_2 &= 10 (0.6438 \times 1 + j1.189 \times 0) = 6.438 \\
 p_3 &= 10 \left( 0.6438 \times \frac{1}{2} - j1.189 \times \frac{\sqrt{3}}{2} \right) = 3.219 - j10.30 \\
 p_4 &= 10 \left( -0.6438 \times \frac{1}{2} - j1.189 \times \frac{\sqrt{3}}{2} \right) = -3.219 - j10.30 \\
 p_5 &= 10 (-0.6438 \times 0 - j1.189 \times 0) = -6.438 \\
 p_6 &= 10 \left( -0.6438 \times \frac{1}{2} + j1.189 \times \frac{\sqrt{3}}{2} \right) = -3.219 + j10.30
 \end{aligned}$$

and the gain adjusted transfer function of the resulting Type 1 filter is

$$H(s) = \frac{750}{(s^2 + 6.438s + 116.5)(s + 6.438)}$$

The pole-zero plot for the Chebyshev Type 1 filter is shown below.



## 2.2 The Chebyshev Type 2 Filter

The Chebyshev Type 2 filter has a monotonically decreasing magnitude function in the pass-band, but introduces equi-amplitude ripple in the stop-band by the inclusion of system zeros on the imaginary axis. The Type 2 filter is defined by the power gain function:

$$|H(j\Omega)|^2 = \frac{1}{1 + \epsilon^2 \frac{T_N^2(\Omega_r/\Omega_c)}{T_N^2(\Omega_r/\Omega)}} \quad (12)$$

If we make the substitutions

$$\nu = \frac{\Omega_r \Omega_c}{\Omega} \quad \text{and} \quad \hat{\epsilon} = \frac{1}{\epsilon T_N(\Omega_r/\Omega_c)}$$

Eq. 24 may be written in terms of the modified frequency  $\nu$

$$|H(j\nu)|^2 = \frac{\hat{\epsilon}^2 T_N^2(\nu/\Omega_c)}{1 + \hat{\epsilon}^2 T_N^2(\nu/\Omega_c)} \quad (13)$$

which has a denominator similar to the Type 1 filter, but has a numerator that contains a Chebyshev polynomial, and is of order  $2N$ . We can use a method similar to that used in the Type 1 filter design to find the poles as follows:

1. First define a complex variable, say  $\tau = \mu + j\nu$  (analogous to the Laplace variable  $s = \sigma + j\Omega$  used in the type 1 design) and write the power transfer function:

$$|H(\tau)|^2 = \frac{\hat{\epsilon}^2 T_N^2(\tau/j\Omega_c)}{1 + \hat{\epsilon}^2 T_N^2(\tau/j\Omega_c)}$$



The poles are found using the method developed for the Type 1 filter, the zeros are found as the roots of the polynomial  $T_N(\tau/j\Omega_c)$  on the imaginary axis  $\tau = j\nu$ . From the definition  $T_N(x) = \cos(N \cos^{-1}(x))$  it is easy to see that the roots of the Chebyshev polynomial occur at

$$x = \cos\left(\frac{(n-1/2)\pi}{N}\right) \quad n = 1 \dots N$$

and from Eq. (25) the system zeros will be at

$$\tau_n = j\Omega_c \cos\left(\frac{(n-1/2)\pi}{N}\right) \quad n = 1 \dots N.$$

2. The poles and zeros are mapped back to the  $s$ -plane using  $s = \Omega_r\Omega_c/\tau$  and the  $N$  left half-plane poles are selected as the poles of the filter.
3. The transfer function is formed and the system gain is adjusted to unity at  $\Omega = 0$ .

## ■ Example 2

Repeat the previous Chebyshev Type 1 design example using a Chebyshev Type 2 filter.

From the previous example we have  $\Omega_c = 10$  rad/s.,  $\Omega_r = 20$  rad/s.,  $\epsilon = 1/3$ ,  $\lambda = 4.358$ . The procedure to find the required order is the same as before, and we conclude that  $N = 3$ . Next, define

$$\begin{aligned} \nu &= \frac{\Omega_r\Omega_c}{\Omega} = \frac{200}{\Omega} \\ \hat{\epsilon} &= \frac{1}{\epsilon T_N(\Omega_r/\Omega_c)} = \frac{3}{T_3(2)} = 0.1154 \end{aligned}$$

Determine  $\alpha$ :

$$\alpha = \frac{1}{N} \sinh^{-1}\left(\frac{1}{\hat{\epsilon}}\right) = \frac{1}{3} \sinh^{-1}(8.666) = 0.9520$$

and  $\sinh \alpha = 1.1024$ , and  $\cosh \alpha = 1.4884$ .

The values of  $\gamma_n = (2n-1)\pi/6$  for  $n = 1 \dots 6$  are the same as the design for the

Type 1 filter, so that the poles of  $|H(\tau)|^2$  are

$$\begin{aligned}
 p_n &= \Omega_c (\sinh \alpha \sin \gamma_n + j \cosh \alpha \cos \gamma_n) \\
 \tau_1 &= 10 \left( 1.1024 \times \frac{1}{2} + j1.4884 \times \frac{\sqrt{3}}{2} \right) = 5.512 + j12.890 \\
 \tau_2 &= 10 (1.1024 \times 1 + j1.4884 \times 0) = 11.024 \\
 \tau_3 &= 10 \left( 1.1024 \times \frac{1}{2} - j1.488 \times \frac{\sqrt{3}}{2} \right) = 5.512 - j12.890 \\
 \tau_4 &= 10 \left( -1.1024 \times \frac{1}{2} - j1.4884 \times \frac{\sqrt{3}}{2} \right) = -5.512 - j12.890 \\
 \tau_5 &= 10 \left( -1.1024 \times \frac{1}{2} - j1.488 \times 0 \right) = -11.024 \\
 \tau_6 &= 10 \left( -1.1024 \times \frac{1}{2} + j1.4884 \times \frac{\sqrt{3}}{2} \right) = -5.512 + j12.890
 \end{aligned}$$

The three left half-plane poles ( $\tau_4, \tau_5, \tau_6$ ) are mapped back to the  $s$ -plane using  $s = \Omega_r \Omega_c / \tau$  giving three filter poles

$$\begin{aligned}
 p_1, p_2 &= -5.609 \pm j13.117 \\
 p_3 &= -18.14
 \end{aligned}$$

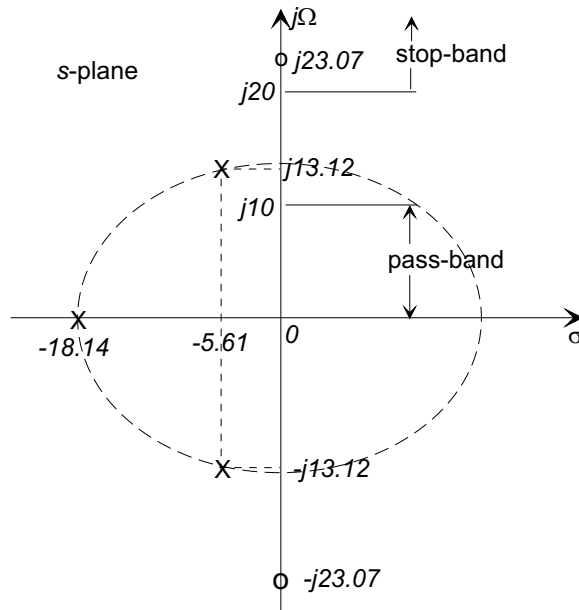
The system zeros are the roots of

$$T_3(\nu/j\Omega_c) = 4(\nu/j\Omega_c)^3 - 3(\nu/j\Omega_c) = 0$$

from the definition of  $T_N(x)$ , giving  $\nu_1 = 0$  and  $\nu_2, \nu_3 = \pm j8.666$ . Mapping these back to the  $s$ -plane gives two finite zeros  $z_1, z_2 = \pm j23.07$ ,  $z_3 = \infty$  (the zero at  $\infty$  does not affect the system response) and the unity gain transfer function is

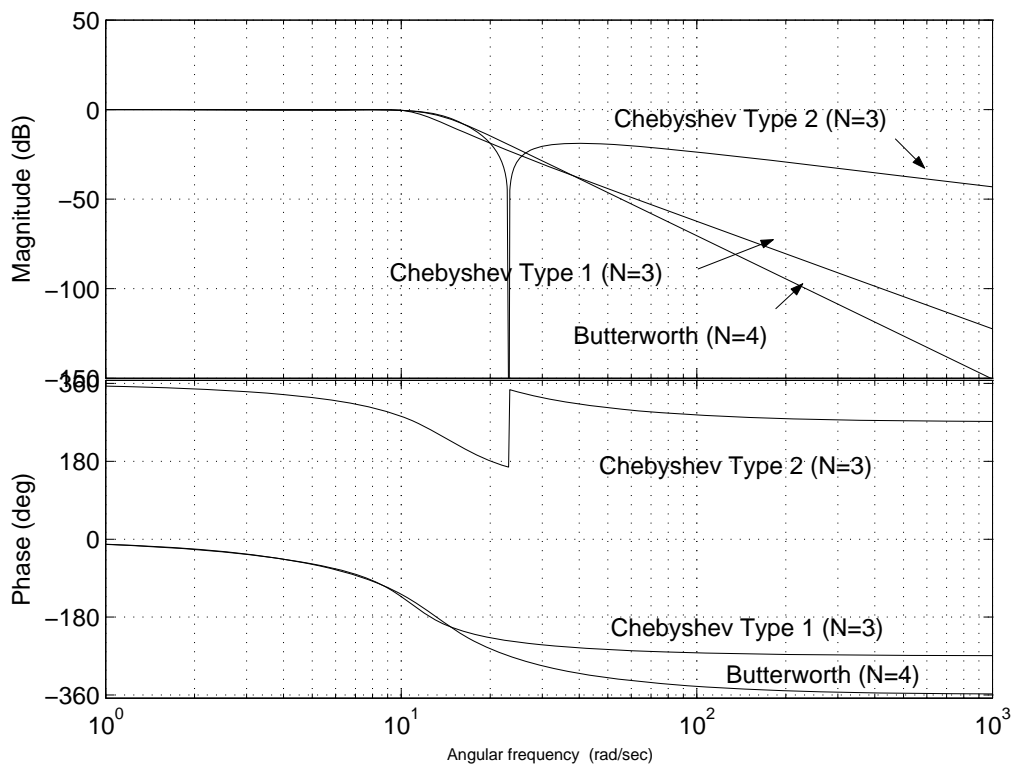
$$\begin{aligned}
 H(s) &= \frac{-p_1 p_2 p_3}{z_1 z_2} \frac{(s - z_1)(s - z_2)}{(s - p_1)(s - p_2)(s - p_3)} \\
 &= \frac{6.9365(s^2 + 532.2)}{(s + 18.14)(s^2 + 11.22s + 203.5)}
 \end{aligned}$$

The pole-zero plot for this filter is shown in below. Note that the poles again lie on ellipse, and the presence of the zeros in the stop-band.



### 2.3 Comparison of Filter Responses

Bode plot responses for the three previous example filters are shown below:



While all filters meet the design specification, it can be seen that the Butterworth and the Chebyshev Type 1 filters are all-pole designs and have an asymptotic high-frequency magnitude slope of  $-20N$  dB/decade, in this case  $-80$  dB/decade for the Butterworth design and  $-60$  dB/decade for the Chebyshev Type 1 design. The Type 2 Chebyshev design has two finite zeros on the imaginary axis at a frequency of  $23.07$  rad/s, forcing the response to zero at this frequency, but with the result that its asymptotic high frequency response has a slope of only  $-20$  dB/decade. Note also the singularity in the phase response of the Type 2 Chebyshev filter, caused by the two purely imaginary zeros.

The pass-band and stop-band power responses are shown in below. Notice that the design method developed here guarantees that the response will meet the specification at the cut-off frequency (in this case  $|H(j\Omega)|^2 = 0.9$  at  $\Omega_c = 10$ ). Other design methods (such as used by MATLAB) may not use this criterion.

