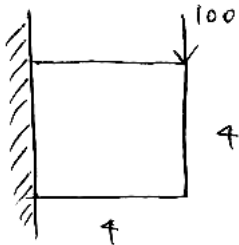


Lecture 6 - Finite element formulation, example, convergence

6.1 Example



$t = 0.1, E, \nu$  plane stress

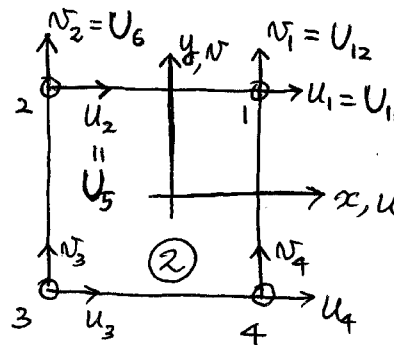
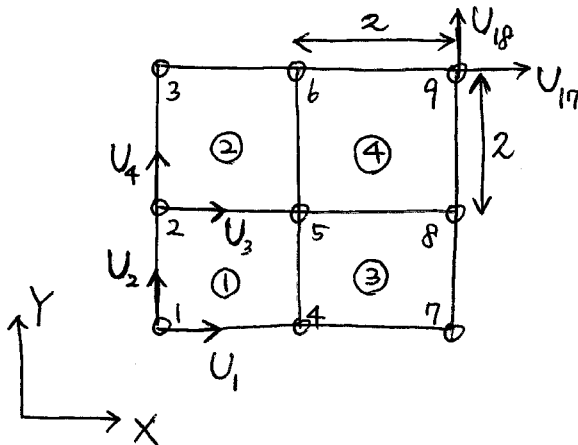
Reading:  
Ex. 4.6 in  
the text

$$KU = R; \quad R = R_B + R_s + R_c + R_r \tag{6.1}$$

$$K = \sum_m K^{(m)}; \quad K^{(m)} = \int_{V^{(m)}} B^{(m)T} C^{(m)} B^{(m)} dV^{(m)} \tag{6.2}$$

$$R_B = \sum_m R_B^{(m)}; \quad R_B^{(m)} = \int_{V^{(m)}} H^{(m)T} f^{B(m)} dV^{(m)} \tag{6.3}$$

6.1.1 F.E. model



$$K|_{\text{el. (2)}} = \begin{bmatrix} \circ & \square & \triangle & \times & \times & \times & \times & \times \\ \vdots & & & & & & & \end{bmatrix} \begin{matrix} \leftarrow u_1 \\ \vdots \end{matrix} \tag{6.4}$$



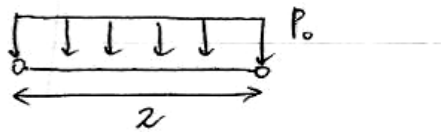
$$\mathbf{H}^S = \mathbf{H} \Big|_{y=+1} \quad (6.10)$$

$$= \begin{bmatrix} \frac{1}{2}(1+x) & \frac{1}{2}(1-x) & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1}{2}(1+x) & \frac{1}{2}(1-x) & 0 & 0 \end{bmatrix} \quad (6.11)$$

From (6.7);

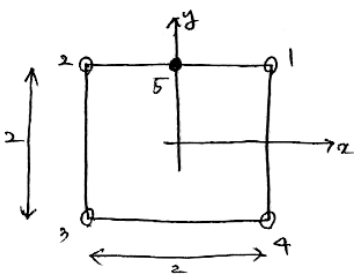
$$\mathbf{R}_S = \int_{-1}^{+1} \begin{bmatrix} \frac{1}{2}(1+x) \\ \frac{1}{2}(1-x) \\ 0 \\ 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ \frac{1}{2}(1+x) \\ \frac{1}{2}(1-x) \end{bmatrix} \begin{bmatrix} 0 \\ -p(x) \end{bmatrix} \underbrace{(0.1)}_{\text{thickness}} dx \quad (6.12)$$

$$\mathbf{R}_S = \begin{bmatrix} 0 \\ 0 \\ -p_0(0.1) \\ -p_0(0.1) \end{bmatrix} \quad (6.13)$$



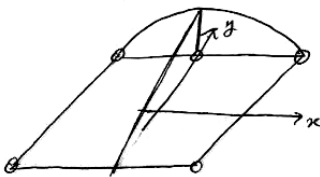
$$\text{total load} = P_0 \times 0.1 \times 2$$

### 6.1.2 Higher-order elements



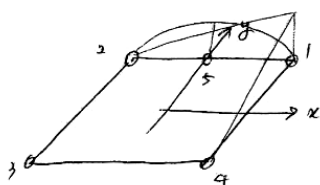
Want  $h_1, h_2, h_3, h_4, h_5$

$$u(x, y) = \sum_{i=1}^5 h_i u_i.$$



$h_i = 1$  at node  $i$  and 0 at all other nodes.

$$h_5 = \frac{1}{2}(1-x^2)(1+y)$$



$$h_1 = \frac{1}{4}(1+x)(1+y) - \frac{1}{2}h_5 \quad (6.14)$$

$$h_2 = \frac{1}{4}(1-x)(1+y) - \frac{1}{2}h_5 \quad (6.15)$$

$$h_3 = \frac{1}{4}(1-x)(1-y) \quad (6.16)$$

$$h_4 = \frac{1}{4}(1+x)(1-y) \quad (6.17)$$

**Note:**

$$\boxed{\sum h_i = 1}$$

We must have  $\sum_i h_i = 1$  to satisfy the rigid body mode condition.

$$u(x, y) = \sum_i h_i u_i \quad (6.18)$$

Assume all nodal point displacements =  $u^*$ . Then,

$$u(x, y) = \sum_i h_i u^* = u^* \sum_i h_i = u^* \quad (6.19)$$

From (6.1),

$$\left( \sum_m \mathbf{K}^{(m)} \right) \mathbf{U} = \mathbf{R} \quad (6.20)$$

$$\sum_m \left[ \int_{V^{(m)}} \mathbf{B}^{(m)T} \mathbf{C}^{(m)} \mathbf{B}^{(m)} dV^{(m)} \right] \mathbf{U} = \mathbf{R} \quad (6.21)$$

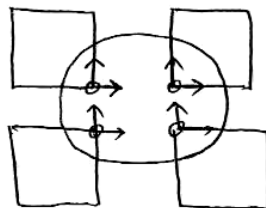
where  $\mathbf{C}^{(m)} \mathbf{B}^{(m)} \mathbf{U} = \boldsymbol{\tau}^{(m)}$ . (Assume we calculated  $\mathbf{U}$ .)

$$\sum_m \int_{V^{(m)}} \mathbf{B}^{(m)T} \boldsymbol{\tau}^{(m)} dV^{(m)} = \mathbf{R} \quad (6.22)$$

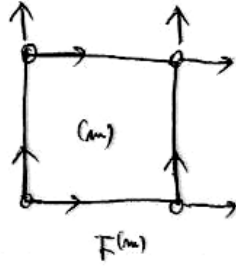
$$\sum_m \mathbf{F}^{(m)} = \mathbf{R}; \quad \mathbf{F}^{(m)} = \int_{V^{(m)}} \mathbf{B}^{(m)T} \boldsymbol{\tau}^{(m)} dV^{(m)} \quad (6.23)$$

### Two properties

I. The sum of the  $\mathbf{F}^{(m)}$ 's at any node is equal to the applied external forces.



II. Every element is in equilibrium under its  $\mathbf{F}^{(m)}$



$$\hat{\mathbf{U}}^T \mathbf{F}^{(m)} = \hat{\mathbf{U}}^T \int_{V^{(m)}} \underbrace{\mathbf{B}^{(m)T} \boldsymbol{\tau}^{(m)}}_{=\bar{\boldsymbol{\epsilon}}^{(m)T}} dV^{(m)} \quad (6.24)$$

$$= \int_{V^{(m)}} \bar{\boldsymbol{\epsilon}}^{(m)T} \boldsymbol{\tau}^{(m)} dV^{(m)} \quad (6.25)$$

$$= 0 \quad (6.26)$$

where  $\hat{\mathbf{U}}^T$  = virtual nodal point displacement.

Apply rigid body displacement.

If we move the element virtually in the rigid body modes,  $\bar{\boldsymbol{\epsilon}}^{(m)}$  is zero. Therefore the virtual work obtained due to virtual motion of the element is zero. Then the element is in equilibrium under its  $\mathbf{F}^{(m)}$ .

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