

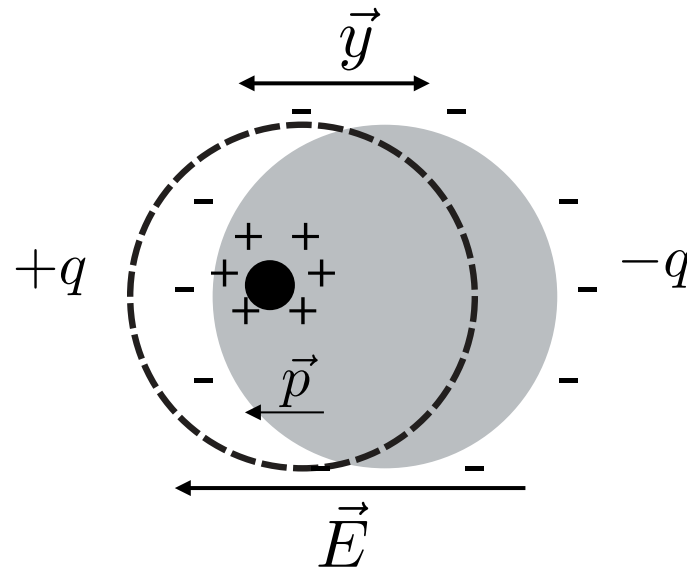
# *Interaction of Atoms and Electromagnetic Waves*

## Outline

- Review: Polarization and Dipoles
- Lorentz Oscillator Model of an Atom
- Dielectric constant and Refractive index

# True or False?

1. The dipole moment of this atom is  $\vec{p} = q\vec{y}$ , and points in the same direction as the polarizing electric field.



2. The susceptibility relates the electric field to the polarization in this form:  $\vec{P} = \epsilon_0 \chi_e \vec{E}$

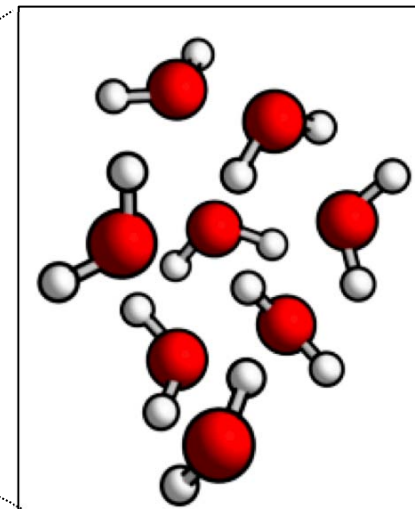
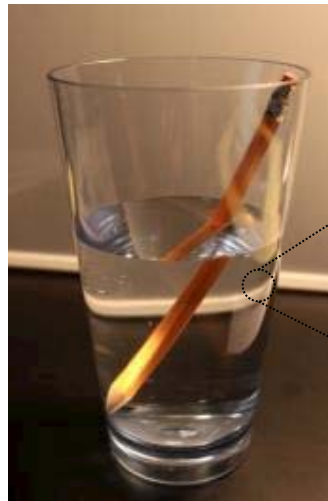
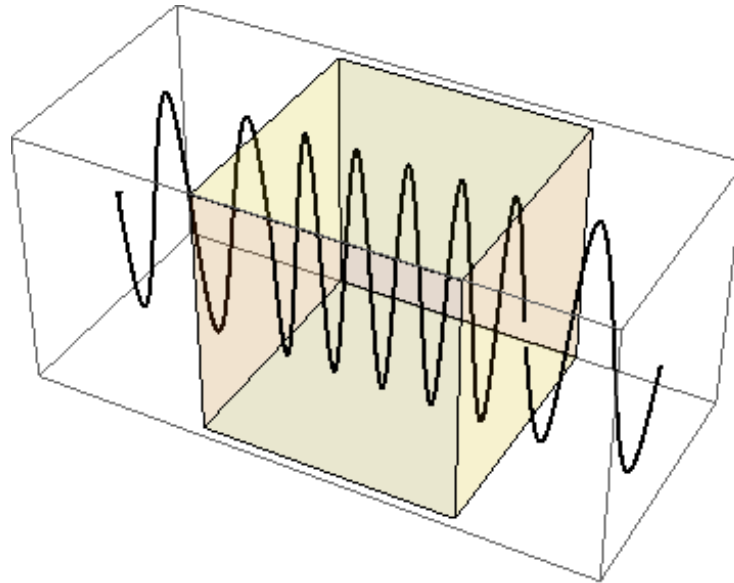
3. The refractive index can be written  $n = \sqrt{\frac{\epsilon_0 \mu_0}{\epsilon \mu}}$

## Refractive Index: Waves in Materials

$$v_p = \frac{1}{\sqrt{\mu\epsilon}}$$

Index of refraction

$$n \equiv \frac{c}{v_p}$$



How do we get from molecules/charges and fields to index of refraction ?

# Index of Refraction

$$\nu \lambda = v_p = c/n$$

↑ frequency      ↑ wavelength

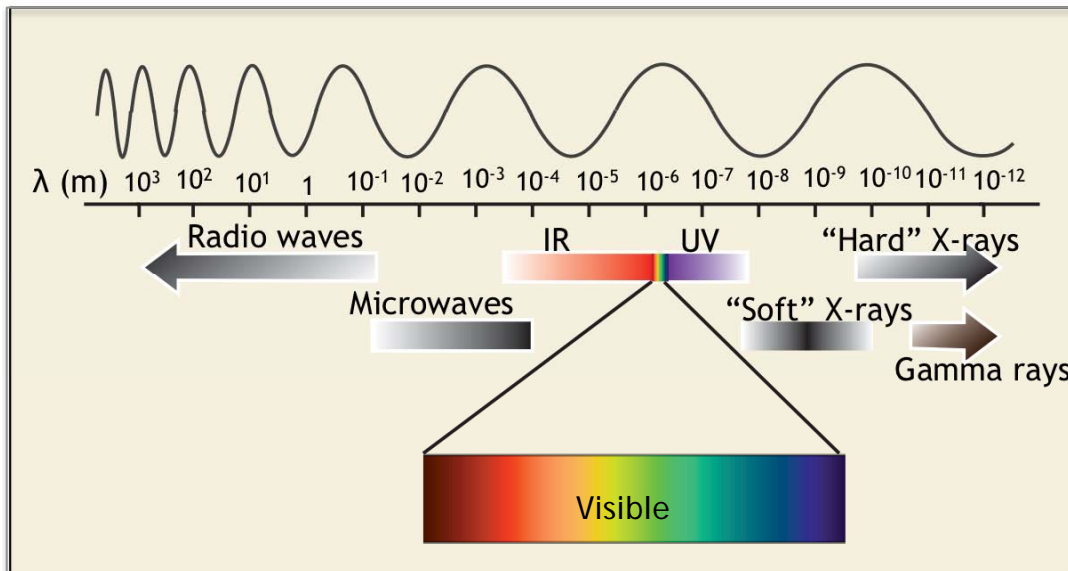
When propagating in a material,

$$c \rightarrow c/n$$

$$\lambda \rightarrow \lambda_o/n$$

$$k \rightarrow k_o/n$$

Material	$n$
Vacuum	1
Air	1.000277
Water liquid	1.3330
Water ice	1.31
Diamond	2.419
Silicon	3.96
at $5 \times 10^{14}$ Hz	



$$E(t, z) = \text{Re}\{\tilde{E}_0 e^{j(\omega t - k_0 n z)}\}$$

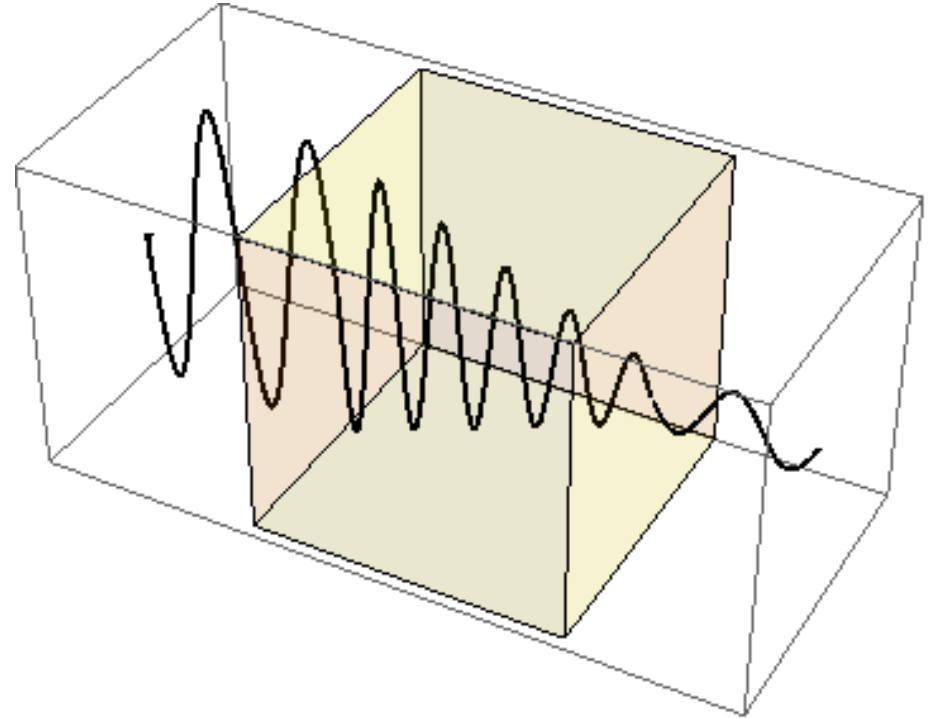


$$E(t, z) = \text{Re}\{\tilde{E}_0 e^{j(\omega t - k z)}\}$$



Photograph by [Hey Paul](#) on Flickr.

## Absorption



Why are these stained glass different colors?

Tomorrow: lump refractive index and absorption into a complex refractive index  $\tilde{n}$

$$E(t, z) = \text{Re}\{ \tilde{E}_0 e^{-\alpha z/2} e^{j(\omega t - k_0 n z)} \}$$

↑
↑

Absorption coefficient                      Refractive index

## Incident Solar Radiation

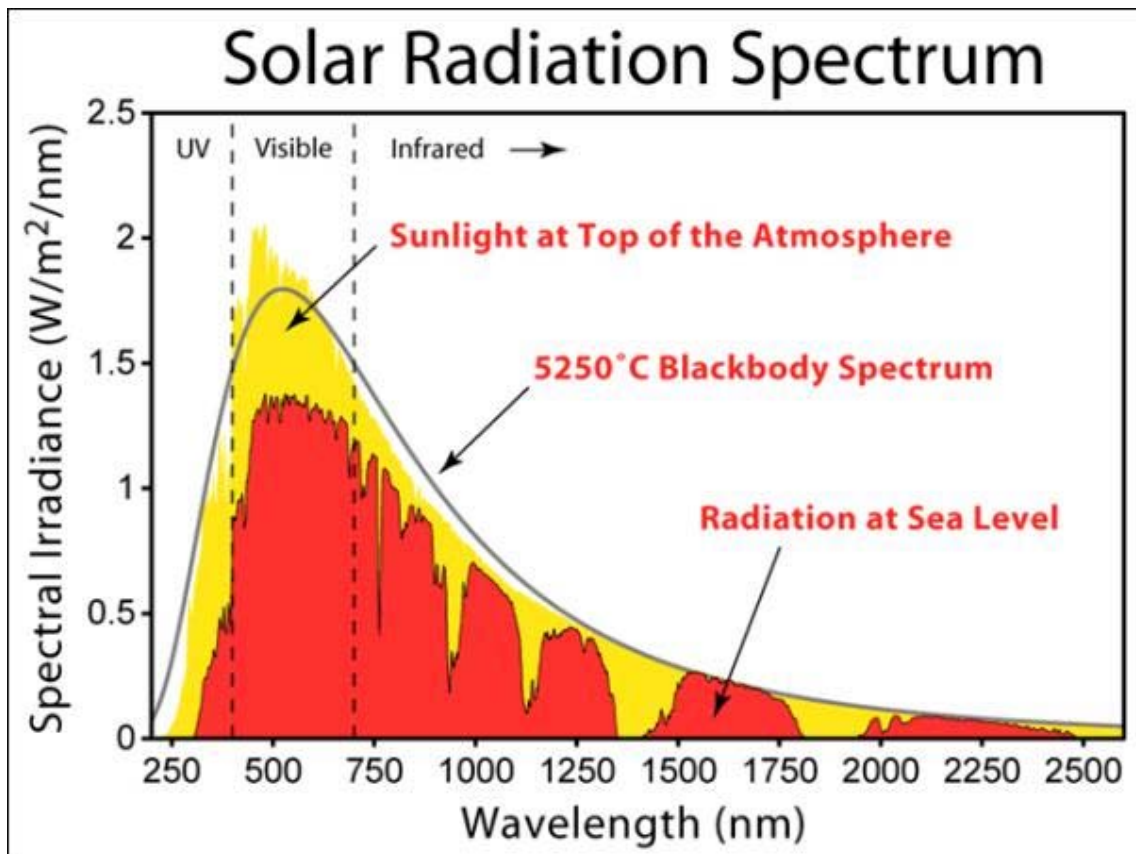
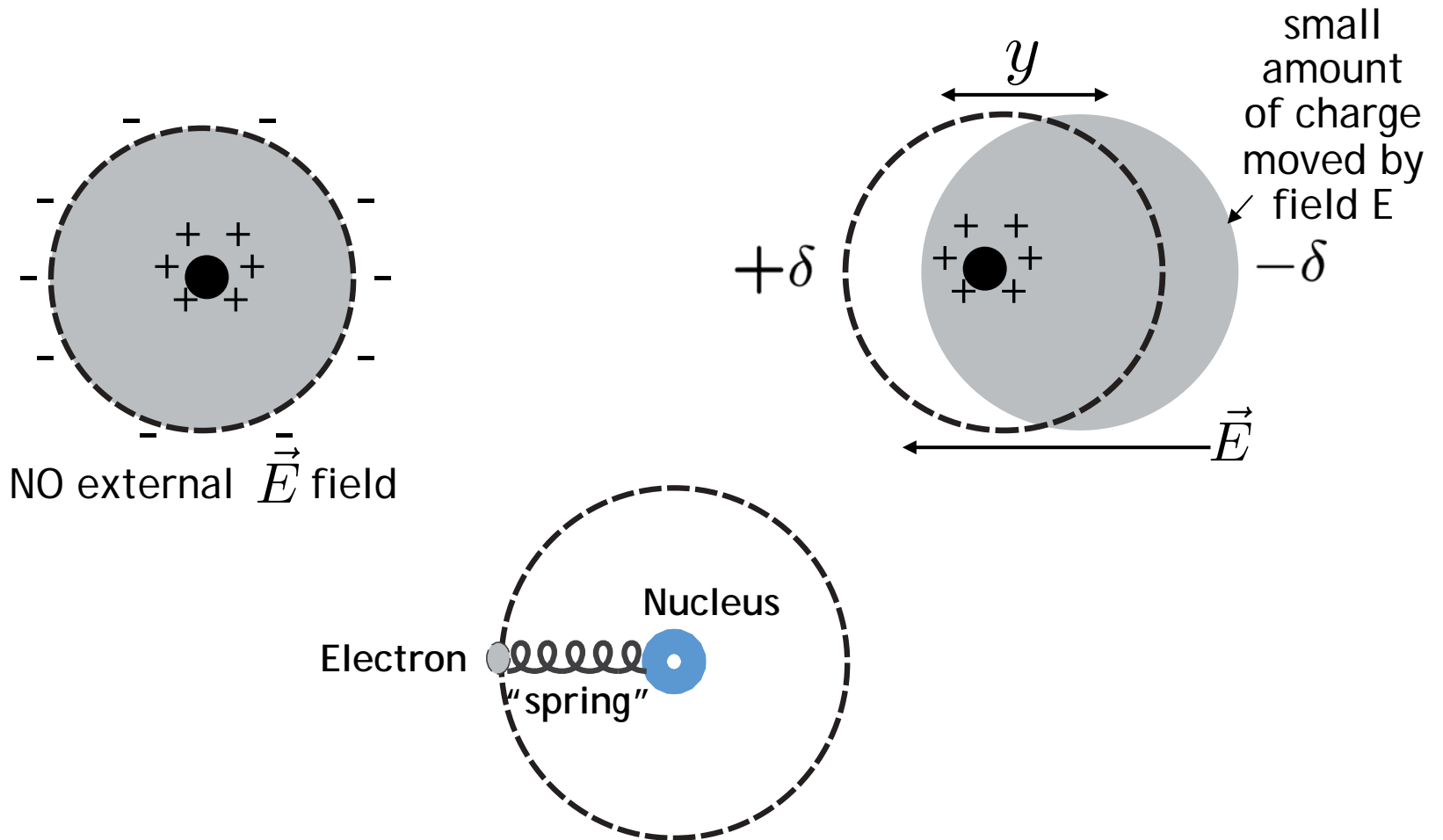


Image created by Robert A. Rohde / Global Warming Art. Used with permission.

How do we introduce propagation through a medium (atmosphere) into Maxwell's Equations?

# Microscopic Description of Dielectric Constant



Density of dipoles...

$$\vec{P} = N\delta\vec{x}$$

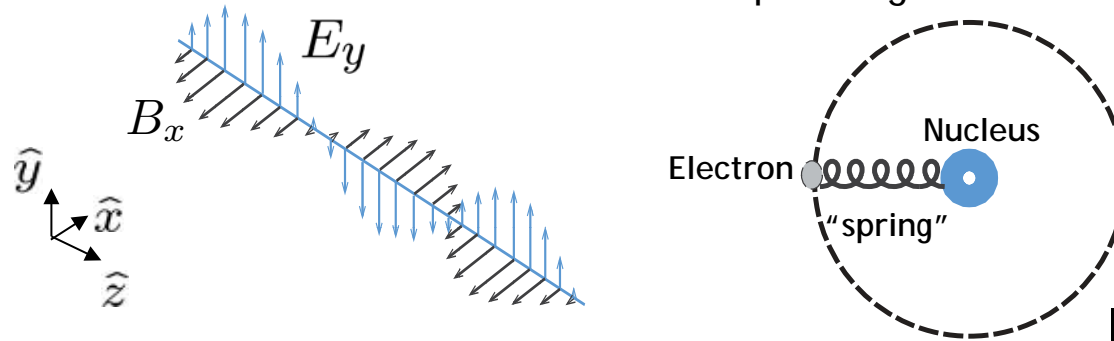
... equivalent to ...

Electric field polarizes molecules...

$$\vec{P} = \epsilon_0\chi_e\vec{E}$$

# Lorentz Oscillator

Lorentz was a late nineteenth century physicist, and quantum mechanics had not yet been discovered. However, he did understand the results of classical mechanics and electromagnetic theory. Therefore, he described the problem of atom-field interactions in these terms. Lorentz thought of an atom as a mass ( the nucleus ) connected to another smaller mass ( the electron ) by a spring. The spring would be set into motion by an electric field interacting with the charge of the electron. The field would either repel or attract the electron which would result in either compressing or stretching the spring.



Hendrik Lorentz (1853-1928)  
Nobel in 1902 for Zeeman Effect

Lorentz was not positing the existence of a physical spring connecting the electron to an atom; however, he did postulate that the force binding the two could be described by Hooke's Law:

$$F(y) = -k_s y$$

where  $y$  is the displacement from equilibrium. If Lorentz's system comes into contact with an electric field, then the electron will simply be displaced from equilibrium. The oscillating electric field of the electromagnetic wave will set the electron into harmonic motion. The effect of the magnetic field can be omitted because it is miniscule compared to the electric field.

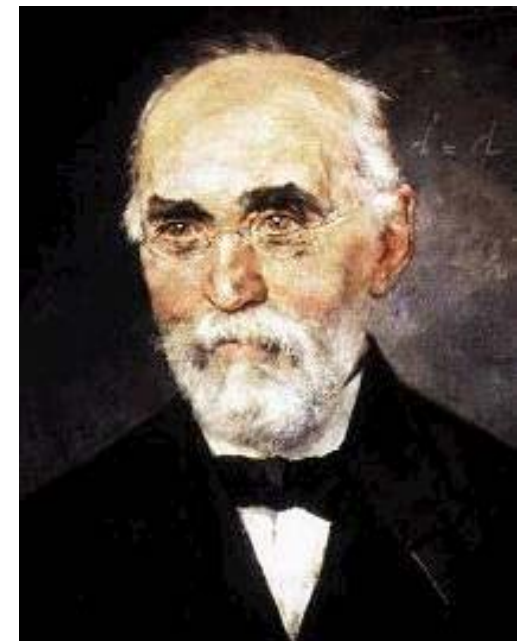


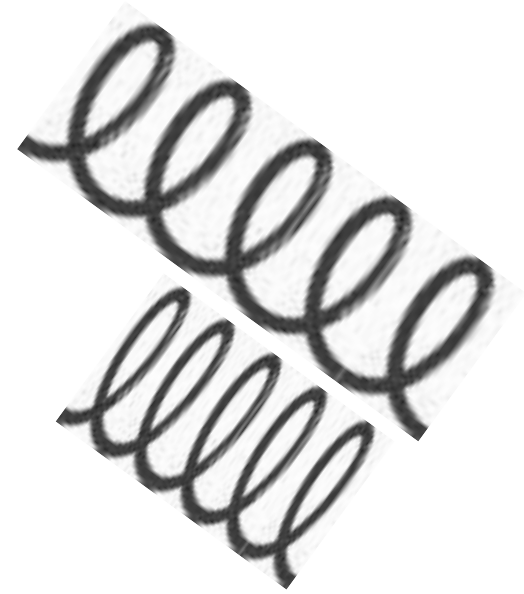
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# Springs have a resonant frequency

Hooke's Law  $m \frac{d^2 y}{dt^2} = -ky$

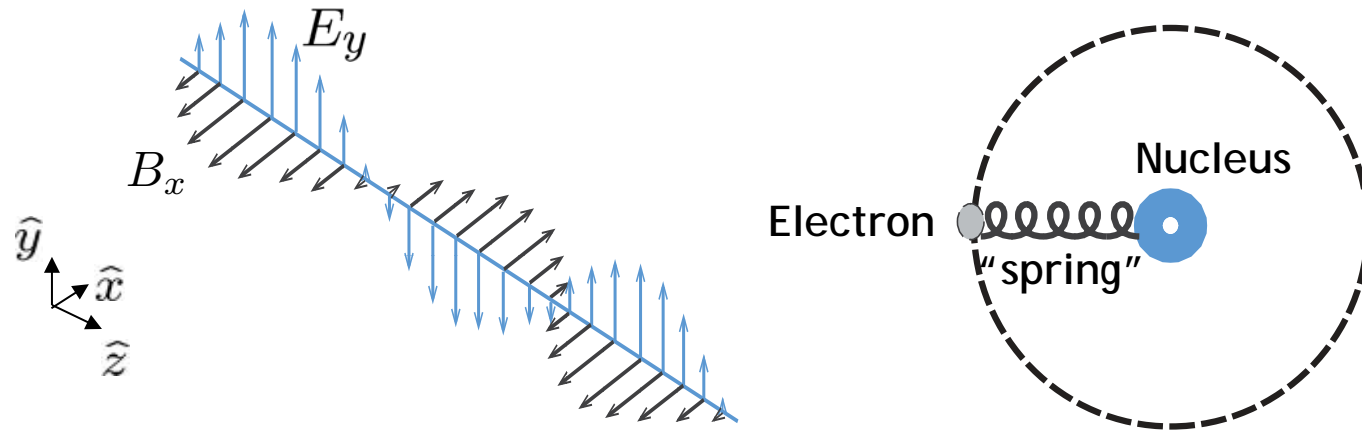
$$\frac{d^2 y}{dt^2} = -\frac{k}{m}y$$



Solution  $y(t) = A \sin\left(t \underbrace{\sqrt{\frac{k}{m}}}_{\omega_0}\right) + B \cos\left(t \underbrace{\sqrt{\frac{k}{m}}}_{\omega_0}\right)$

So we can write:  $\omega_0^2 = \frac{k}{m} \quad \rightarrow \quad k = m\omega_0^2$

# Microscopic Description of Dielectric Constant



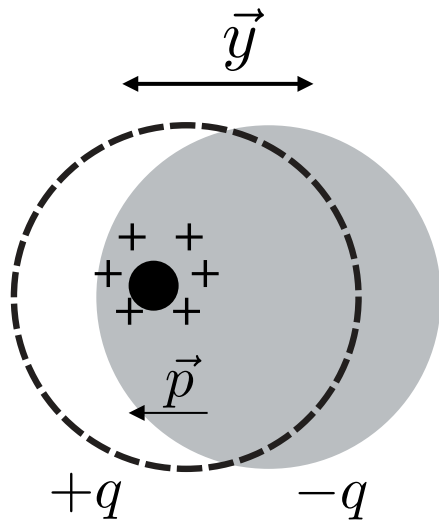
$$m \frac{d^2 y}{dt^2} = -m\omega_o^2 y + qE_y - m\gamma \frac{dy}{dt}$$

Electron mass  $\nearrow$   $m \frac{d^2 y}{dt^2}$   $\nwarrow$  Restoring force (binding electron & nucleus)  $\nwarrow$   $qE_y$   $\nwarrow$   $\vec{E}$  field force  $\nwarrow$   $-m\gamma \frac{dy}{dt}$   $\nwarrow$  Damping

## Solution using complex variables

Lets plug-in the expressions for  $E_y$  and  $y$  into the differential equation from slide 9:

Natural resonance



$$\frac{d^2}{dt^2}y(t) + \gamma \frac{d}{dt}y(t) + \omega_o^2 y(t) = \frac{q}{m} E_y(t)$$

$$E_y(t) = \text{Re}\{E_y e^{j\omega t}\} \quad y(t) = \text{Re}\{y e^{j\omega t}\}$$

$$\omega^2 y + j\omega\gamma y + \omega_o^2 y = \frac{q}{m} E_y$$

$$y = \frac{q}{m} \frac{1}{(\omega_o^2 - \omega^2) + j\omega\gamma} E_y$$

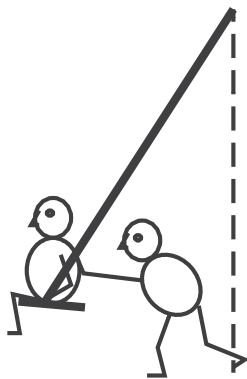
## Oscillator Resonance

$$y = \frac{q}{m} \frac{1}{(\omega_0^2 - \omega^2) + j\omega\gamma} E_y$$

$$E_y(t) = \text{Re}\{E_y e^{j\omega t}\}$$

$$y(t) = \text{Re}\{y e^{j\omega t}\}$$

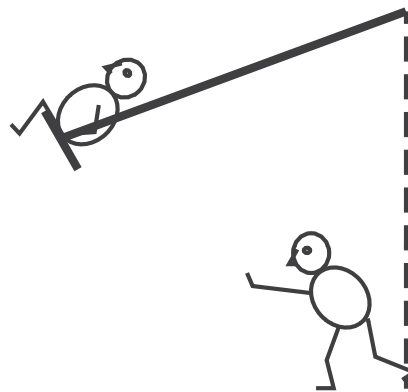
Driven harmonic oscillator: **Amplitude** and **Phase** depend on frequency



**Low** frequency

medium amplitude

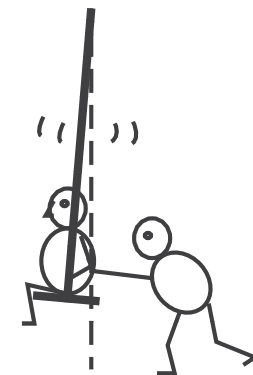
Displacement,  $y$   
in phase with  $E_y$



**At resonance**

large amplitude

Displacement,  $y$   
 $90^\circ$  out of phase with  $E_y$



**High** frequency

vanishing amplitude

Displacement  $y$  and  $E_y$   
in antiphase

## Polarization

Since charge displacement,  $y$ , is directly related to polarization,  $P$ , of our material we can then rewrite the differential equation:

$$\vec{D} = \epsilon_o \vec{E} + \vec{P}$$

For linear polarization in  $\hat{y}$  direction

$$P_y = Nqy$$

$$\left( \frac{d^2}{dt^2} + \gamma \frac{d}{dt} + \omega_o^2 \right) P_y(t) = \frac{Nq^2}{m} E_y(t) = \epsilon_o \omega_p^2 E_y(t)$$

$$\omega_p^2 = \frac{Nq^2}{\epsilon_o m}$$

$$P_y(t) = \text{Re}\{P_y e^{j\omega t}\}$$

$$P_y = \frac{\omega_p^2}{(\omega_o^2 - \omega^2) + j\gamma\omega} \epsilon_o E_y$$

## Dielectric Constant from the Lorentz Model

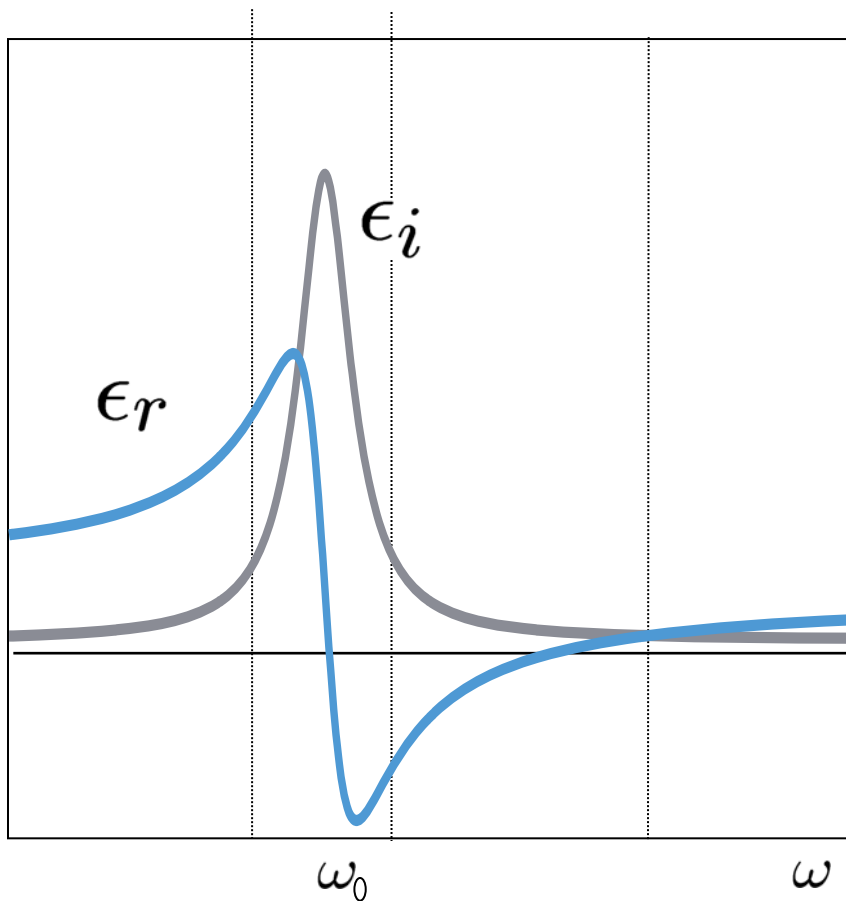
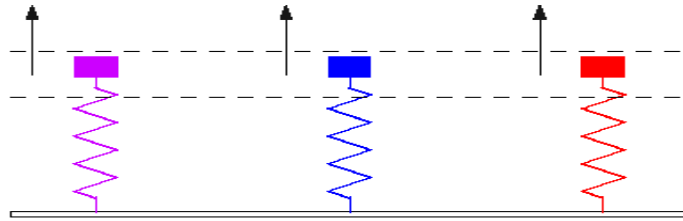
$$\vec{D} = \epsilon_o \vec{E} + \vec{P}$$

$$\vec{P} = \frac{\omega_p^2}{(\omega_o^2 - \omega^2) + j\gamma\omega} \epsilon_o \vec{E}$$

$$\vec{D} = \epsilon_o \left[ 1 + \frac{\omega_p^2}{(\omega_o^2 - \omega^2) + j\gamma\omega} \right] \vec{E}$$

$$\epsilon = \epsilon_o \left[ 1 + \frac{\omega_p^2}{(\omega_o^2 - \omega^2) + j\gamma\omega} \right]$$

## Microscopic Lorentz Oscillator Model



$$\vec{P} = \frac{\omega_p^2}{(\omega_0^2 - \omega^2) + j\gamma\omega} \epsilon_0 \vec{E}$$

$$\omega_p^2 = \frac{Nq^2}{\epsilon_0 m}$$

$$\epsilon = \epsilon_r - j\epsilon_i$$

## Real and imaginary parts

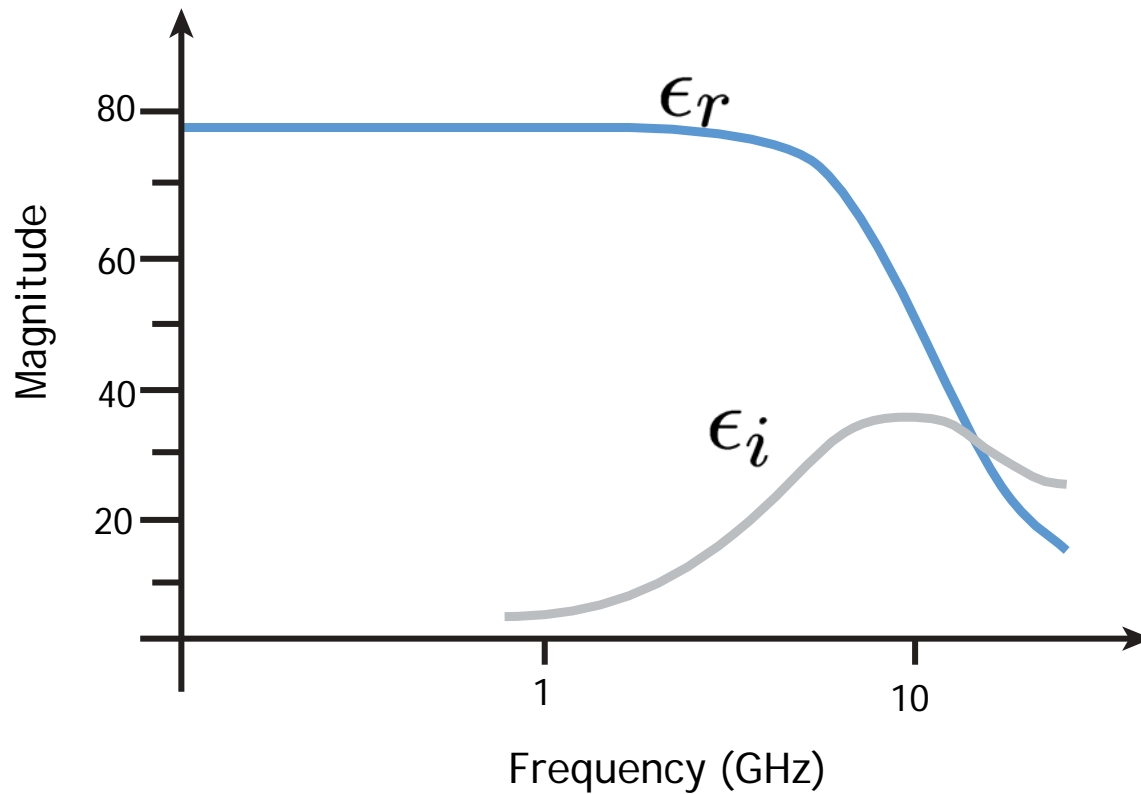
$$\begin{aligned}\epsilon &= \frac{\omega_p^2}{(\omega_o^2 - \omega^2) + j\omega\gamma} \\ &= \frac{\omega_p^2(\omega_o^2 - \omega^2)}{(\omega_o^2 - \omega^2)^2 + \omega^2\gamma^2} - j \frac{\omega_p^2\omega\gamma}{(\omega_o^2 - \omega^2)^2 + \omega^2\gamma^2}\end{aligned}$$

$$\epsilon = \epsilon_r - j\epsilon_i$$

$$\epsilon_r = \frac{\omega_p^2(\omega_o^2 - \omega^2)}{(\omega_o^2 - \omega^2)^2 + \omega^2\gamma^2} \quad \epsilon_i = \frac{\omega_p^2\omega\gamma}{(\omega_o^2 - \omega^2)^2 + \omega^2\gamma^2}$$

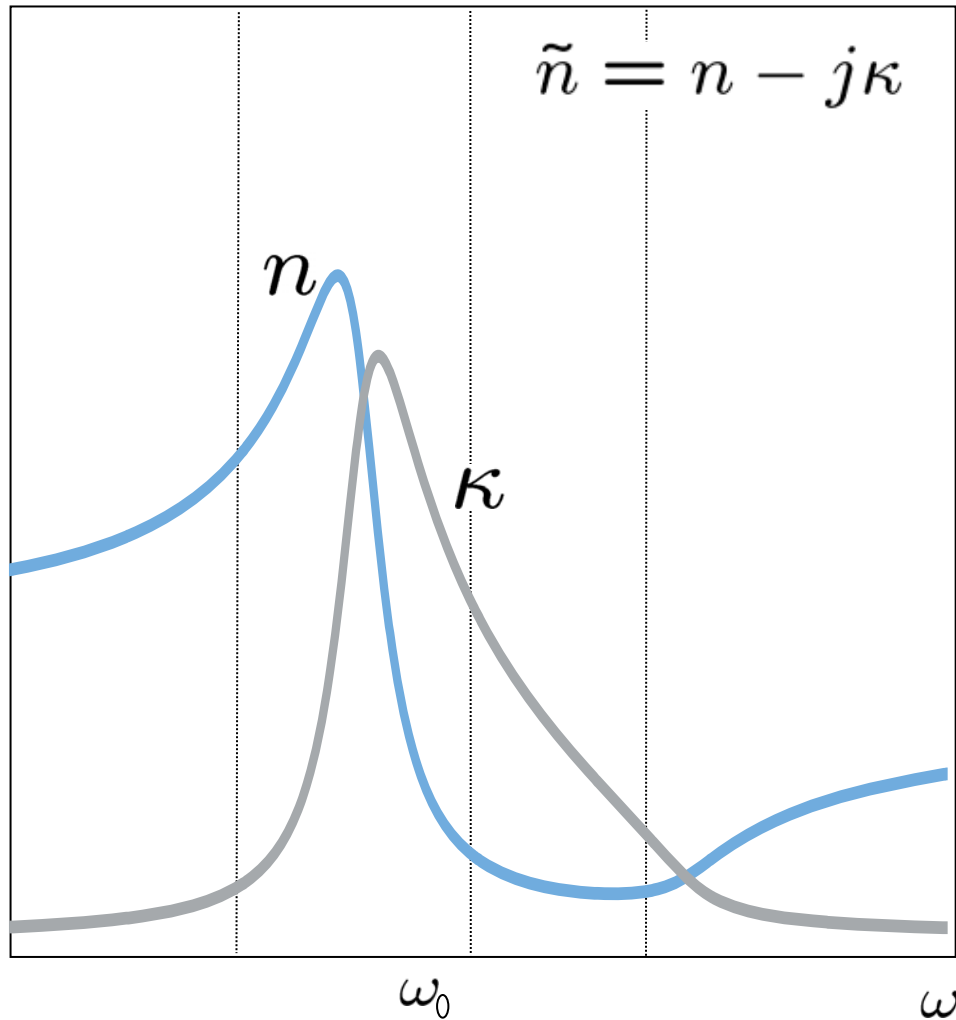


## Dielectric constant of water



Microwave ovens usually operate at 2.45 GHz

## Complex Refractive Index



$$\tilde{n} = n - j\kappa$$

$$\tilde{\epsilon} = \tilde{n}^2$$

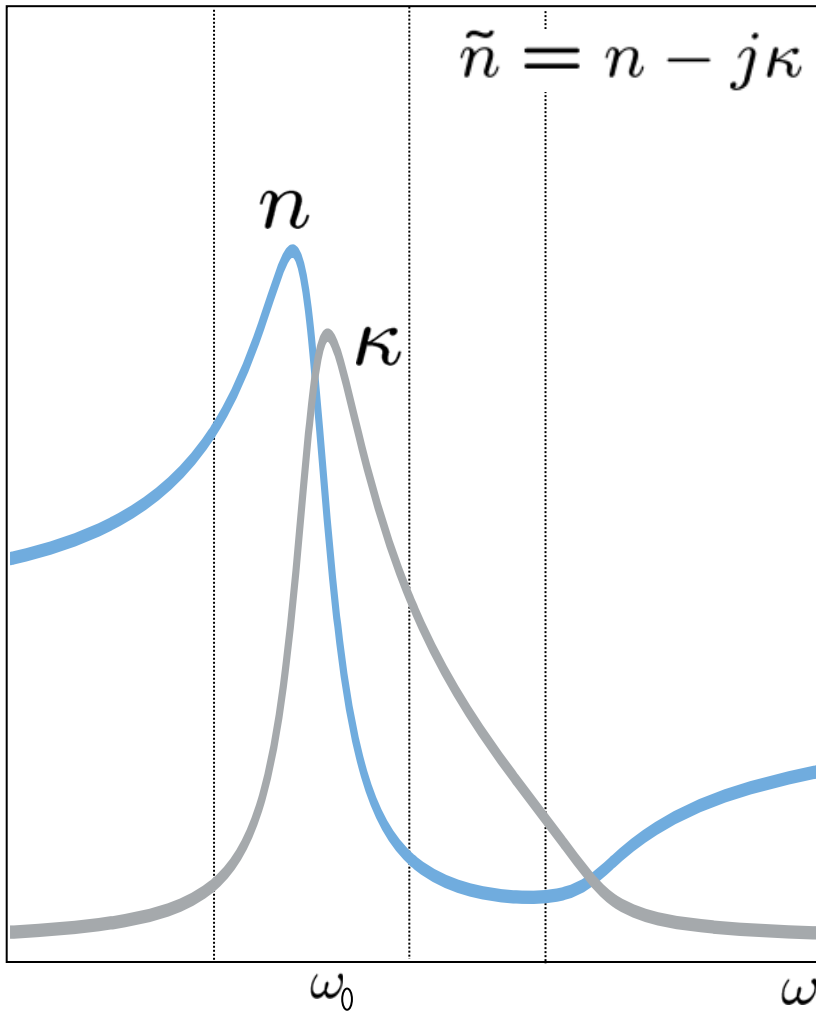
$$= (n - j\kappa)^2$$

$$= n^2 - \kappa^2 - 2jn\kappa$$

$$\epsilon_r = n^2 - \kappa^2$$

$$\epsilon_i = 2n_r n_i$$

## Absorption Coefficient



$$n = \frac{1}{\sqrt{2}} \sqrt{\epsilon_r + \sqrt{\epsilon_r^2 + \epsilon_i^2}}$$

$$\kappa = \frac{1}{\sqrt{2}} \sqrt{-\epsilon_r + \sqrt{\epsilon_r^2 + \epsilon_i^2}}$$

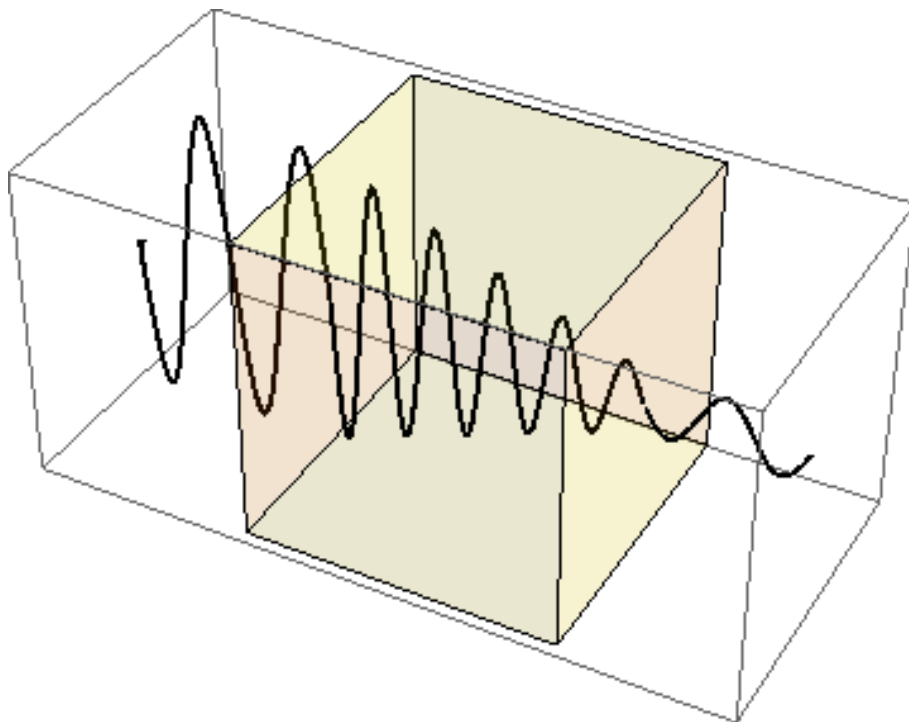
$$E(t, z) = \text{Re}\{ \tilde{E}_0 e^{-\alpha z/2} e^{j(\omega t - k_0 n z)} \}$$

Absorption

Refractive  
index

$$\alpha = 2k_0 \kappa = 2 \frac{2\pi}{\lambda_0} \kappa \quad [\text{cm}^{-1}]$$

## Absorption



$$E(t, z) = \text{Re}\left\{ \tilde{E}_0 e^{-\alpha z/2} e^{j(\omega t - k_0 n z)} \right\}$$

Absorption  
coefficient

Refractive  
index

$$I(z) = I_0 e^{-\alpha z} \quad \text{Beer-Lambert Law or Beer's Law}$$



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