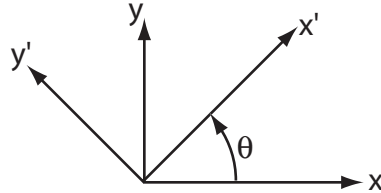


## Coordination Transformations for Strain & Stress Rates

To keep the presentation as simple as possible, we will look at purely two-dimensional stress-strain rates. Given an original coordinate system  $(x, y)$  and a rotated system  $(\hat{x}, \hat{y})$  as shown below:



Recall that the strain rates in the  $x$ - $y$  coordinate system are:

$$\varepsilon_{xx} = \frac{\partial u}{\partial x} \quad \varepsilon_{xy} = \frac{1}{2} \left( \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \quad \varepsilon_{yy} = \frac{\partial v}{\partial y}$$

Or, in index notation:

$$\varepsilon_{ij} = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

Also, we note that the unit vectors for the rotated axes are:

$$\begin{aligned} \hat{i} &= \cos \theta \bar{i} + \sin \theta \bar{j} \\ \hat{j} &= -\sin \theta \bar{i} + \cos \theta \bar{j} \end{aligned}$$

Thus, the location of a point in  $(\hat{x}, \hat{y})$  is:

$$\begin{bmatrix} \hat{x} \\ \hat{y} \end{bmatrix} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

Similarly, the velocity components are related by:

$$\begin{bmatrix} \hat{u} \\ \hat{v} \end{bmatrix} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} u \\ v \end{bmatrix}$$

For differential changes, we also have

$$\begin{bmatrix} d\hat{x} \\ d\hat{y} \end{bmatrix} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} dx \\ dy \end{bmatrix}$$

Thus, defining  $T$  as the rotation matrix, we note that:

$$T = \begin{bmatrix} \frac{\partial \hat{x}}{\partial x} & \frac{\partial \hat{x}}{\partial y} \\ \frac{\partial \hat{y}}{\partial x} & \frac{\partial \hat{y}}{\partial y} \end{bmatrix} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix}$$

Inverting this:

$$T^{-1} = \begin{bmatrix} \frac{\partial x}{\partial \hat{x}} & \frac{\partial x}{\partial \hat{y}} \\ \frac{\partial y}{\partial \hat{x}} & \frac{\partial y}{\partial \hat{y}} \end{bmatrix} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix}^{-1} = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

Thus to find  $\frac{\partial \hat{u}}{\partial \hat{x}}$  in terms of  $u, v$  and their derivatives:

$$\frac{\partial \hat{u}}{\partial \hat{x}} = \frac{\partial \hat{u}}{\partial x} \frac{\partial x}{\partial \hat{x}} + \frac{\partial \hat{u}}{\partial y} \frac{\partial y}{\partial \hat{x}} = \frac{\partial \hat{u}}{\partial x} \cos \theta + \frac{\partial \hat{u}}{\partial y} \sin \theta$$

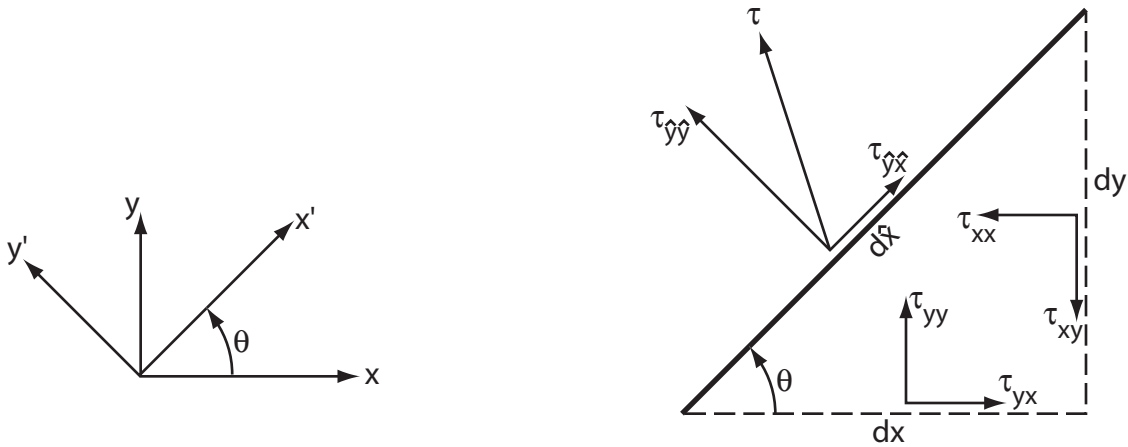
Then, substituting  $\hat{u} = u \cos \theta + v \sin \theta$ :

$$\begin{aligned} \frac{\partial \hat{u}}{\partial \hat{x}} &= \left[ \cos \theta \frac{\partial}{\partial x} + \sin \theta \frac{\partial}{\partial y} \right] (u \cos \theta + v \sin \theta) \\ &= \cos^2 \theta \frac{\partial u}{\partial x} + \cos \theta \sin \theta \frac{\partial v}{\partial x} + \cos \theta \sin \theta \frac{\partial u}{\partial y} + \sin^2 \theta \frac{\partial v}{\partial y} \\ \Rightarrow \boxed{\varepsilon_{\hat{x}\hat{x}} &= \cos^2 \theta \varepsilon_{xx} + 2 \cos \theta \sin \theta \varepsilon_{xy} + \sin^2 \theta \varepsilon_{yy}} \\ \frac{\partial \hat{v}}{\partial \hat{y}} &= \left[ -\sin \theta \frac{\partial}{\partial x} + \cos \theta \frac{\partial}{\partial y} \right] (-u \sin \theta + v \cos \theta) \\ &= \sin^2 \theta \frac{\partial u}{\partial x} - \sin \theta \cos \theta \frac{\partial v}{\partial x} - \sin \theta \cos \theta \frac{\partial u}{\partial y} + \cos^2 \theta \frac{\partial v}{\partial y} \\ \Rightarrow \boxed{\varepsilon_{\hat{y}\hat{y}} &= \sin^2 \theta \varepsilon_{xx} + 2 \sin \theta \cos \theta \varepsilon_{xy} + \cos^2 \theta \varepsilon_{yy}} \\ \frac{1}{2} \left( \frac{\partial \hat{u}}{\partial \hat{y}} + \frac{\partial \hat{v}}{\partial \hat{x}} \right) &= \frac{1}{2} \left\{ \left[ -\sin \theta \frac{\partial}{\partial x} + \cos \theta \frac{\partial}{\partial y} \right] (u \cos \theta + v \sin \theta) + \left[ \cos \theta \frac{\partial}{\partial x} + \sin \theta \frac{\partial}{\partial y} \right] (-\sin \theta + v \cos \theta) \right\} \\ \Rightarrow \boxed{\varepsilon_{\hat{x}\hat{y}} &= \sin \theta \cos \theta (\varepsilon_{yy} - \varepsilon_{xx}) + (\cos^2 \theta - \sin^2 \theta) \varepsilon_{xy}} \end{aligned}$$

If  $x - y$  are the principal strain directions, then  $\varepsilon_{xy} = 0$  and  $\varepsilon_{\hat{x}\hat{x}}, \varepsilon_{\hat{y}\hat{y}}$  &  $\varepsilon_{\hat{y}\hat{x}}$  are

$$\begin{cases} \varepsilon_{\hat{x}\hat{x}} = \cos^2 \theta \varepsilon_{xx} + \sin^2 \theta \varepsilon_{yy} \\ \varepsilon_{\hat{y}\hat{y}} = \sin^2 \theta \varepsilon_{xx} + \cos^2 \theta \varepsilon_{yy} \\ \varepsilon_{\hat{y}\hat{x}} = \sin \theta \cos \theta (\varepsilon_{yy} - \varepsilon_{xx}) \end{cases} \quad \text{if } \varepsilon_{xy} = 0$$

The next step is to relate the stresses in  $(x, y)$  to  $(\hat{x}, \hat{y})$ . Consider a differential surface with  $\hat{y}$  normal:



The resultant stress is given as the vector  $\bar{\tau}$  and the force on the surface is  $\bar{\tau} ds$ . Decomposing the stress vector into the coordinate axes gives:

$$\begin{aligned} \bar{\tau} d\hat{x} &= (\tau_{\hat{y}\hat{x}} \bar{i} + \tau_{\hat{y}\hat{y}} \bar{j}) d\hat{x} \\ &= (-\tau_{xx} dy + \tau_{yx} dx) \bar{i} + (\tau_{yy} dx - \tau_{xy} dy) \bar{j} \end{aligned}$$

Note that:

$$\begin{aligned} dx &= \cos \theta d\hat{x} \\ dy &= \sin \theta d\hat{x} \\ \bar{i} &= \cos \theta \hat{i} - \sin \theta \hat{j} \\ \bar{j} &= \sin \theta \hat{i} + \cos \theta \hat{j} \end{aligned}$$

Thus, the second line becomes:

$$\begin{aligned}
 & (-\tau_{xx} \sin \theta + \tau_{yx} \cos \theta)(\cos \theta \hat{i} - \sin \theta \hat{j}) d\hat{x} \\
 & + (\tau_{yy} \cos \theta - \tau_{xy} \sin \theta)(\sin \theta \hat{i} + \cos \theta \hat{j}) d\hat{x} = (\tau_{\hat{j}\hat{x}} \hat{i} + \tau_{\hat{j}\hat{y}} \hat{j}) d\hat{x}
 \end{aligned}$$

So collecting all the  $\hat{i}$  &  $\hat{j}$  terms (and enforcing  $\tau_{xy} = \tau_{yx}$ ) gives:

$$\begin{aligned}
 \tau_{\hat{j}\hat{x}} &= (\tau_{yy} - \tau_{xx}) \sin \theta \cos \theta + \tau_{yx} (\cos^2 \theta - \sin^2 \theta) \\
 \tau_{\hat{j}\hat{y}} &= \tau_{xx} \sin^2 \theta + \tau_{yy} \cos^2 \theta - 2\tau_{yx} \sin \theta \cos \theta
 \end{aligned}$$

For the principal strain axes,

$$\begin{aligned}
 \tau_{xx} &= 2\mu\varepsilon_{xx} + \lambda(\varepsilon_{xx} + \varepsilon_{yy}) \\
 \tau_{yy} &= 2\mu\varepsilon_{yy} + \lambda(\varepsilon_{xx} + \varepsilon_{yy}) \\
 \tau_{xy} &= 0
 \end{aligned}$$

Plugging this into  $\tau_{\hat{j}\hat{x}}$  and  $\tau_{\hat{j}\hat{y}}$  gives

$$\begin{aligned}
 \tau_{\hat{j}\hat{x}} &= 2\mu \underbrace{(\varepsilon_{yy} - \varepsilon_{xx})}_{\varepsilon_{\hat{y}\hat{x}}} \sin \theta \cos \theta \\
 \tau_{\hat{j}\hat{y}} &= 2\mu \underbrace{(\varepsilon_{xx} \sin^2 \theta + \varepsilon_{yy} \cos^2 \theta)}_{\varepsilon_{\hat{y}\hat{y}}} + \lambda \underbrace{(\varepsilon_{xx} + \varepsilon_{yy})}_{\varepsilon_{\hat{x}\hat{x} + \varepsilon_{\hat{y}\hat{y}}}}
 \end{aligned}$$

Thus, we arrive at the known result:

$$\begin{aligned}
 \tau_{\hat{j}\hat{x}} &= 2\mu\varepsilon_{\hat{x}\hat{y}} \\
 \tau_{\hat{j}\hat{y}} &= 2\mu\varepsilon_{\hat{y}\hat{y}} + \lambda(\varepsilon_{\hat{x}\hat{x}} + \varepsilon_{\hat{y}\hat{y}})
 \end{aligned}$$

A similar derivation would give:

$$\tau_{\hat{x}\hat{x}} = 2\mu\varepsilon_{\hat{x}\hat{x}} + \lambda(\varepsilon_{\hat{x}\hat{x}} + \varepsilon_{\hat{y}\hat{y}})$$